# Scuola Internazionale Superiore di Studi Avanzati 

Geometry and Mathematical Physics

# Introduction to linear ODEs in the Complex Domain and Isomonodromy Deformations 

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## 1 Introduction

solve integrable systems, physics problems such as Schrödinger equation.
Bessel equation

$$
z^{2} y^{\prime \prime}+z y^{\prime}+\left(z^{2}-\nu^{2}\right) y=0
$$

where $\nu \in \mathbb{C}$.
Airy

$$
y^{\prime \prime}=z y
$$

Gauss equation

$$
z(1-z) y^{\prime \prime}+(\gamma-(\alpha+\beta+1) z) y^{\prime}-\alpha \beta y=0
$$

where $\alpha, \beta, \gamma \in \mathbb{C}$.
The Bessel equation occurs in solving the Laplace equation $\nabla \varphi=0$, and the Gauss equation occurs when solving the Schrödinger equation.

More generally, using the notation $y^{(m)}=\frac{d^{m} y}{d z^{m}}$, a linear ODE of order $m$ has the form

$$
y^{(m)}+a_{1}(z) y^{(m-1)}+\cdots+a_{m}(z) y=b(z) .
$$

As we will see, this is a special case of a system of ODEs

$$
\frac{d y}{d z}=\underbrace{\left(\begin{array}{ccc}
a_{11(z)} & \cdots & a_{1 m}(z) \\
\vdots & & \vdots \\
a_{n 1}(z) & \cdots & a_{n m}(z)
\end{array}\right)}_{A(z)} y
$$

and $y$ can be either a vector or a matrix. This is a linear system of order 1 .
What if the matrix is not analytic? That is, if the $a_{i j}$ 's have poles in the complex plane? The question is

- locally, how to represent the solutions closed to singularities?
- globally, how to connect the solutions around different singularities?

Another theme is the study of monodromy propertie: if we take a solution and take a loop around a singularity, the solution will transform under $y(z) \mapsto$ $y(z) M_{\gamma}$, where $M_{\gamma}$ is the monodromy matrix. We will see the result depends only on the homotopy class of the path, and can be described by the monodromy group.

If we have an ODE where the matrix depends on parameters

$$
\frac{d Y}{d z}=A\left(z, t_{1}, \ldots, t_{p}\right) Y
$$

then not only will the solution $Y(z, t)$ depends on the parameters $t$, but also will the singularities and the monomromy matrix $M_{\gamma}(t)$. We will in particular study monodromy "data" that do not depend on $t$. This is called monodromy preserving deformation.

We may also require

$$
\frac{\partial Y}{\partial t_{j}}=\Omega_{j}(z, t) Y, j=1, \ldots, p
$$

subject to the compatibility conditions

$$
\begin{aligned}
\frac{\partial^{2} Y}{\partial t_{j} \partial t_{k}} & =\frac{\partial^{2} Y}{\partial t_{k} \partial t_{j}} \\
\frac{\partial^{2} Y}{\partial z \partial t_{j}} & =\frac{\partial^{2} Y}{\partial t_{j} \partial z}
\end{aligned}
$$

These compatibilities are nonlinear differential equations for $A$ and $\Omega_{j}$ 's.
These compatibilies arise naturally in physics problems, such as

- the Painlevé equations,
- structure of manifolds (Dubrovin-Frobenius manifolds),
- random matrices
- nonlinear PDEs such as KdV and KP


### 1.1 Preliminaries

If $f$ is holomorphic on a domain $B$ then $f$ is smooth. The reason is Cauchy integral formula

$$
f(z)=\frac{1}{2 \pi i} \oint_{\gamma} \frac{f\left(z^{\prime}\right)}{z^{\prime}-z} d z^{\prime}
$$

for $\gamma$ a simple counterclockwise path around $z$, and if $w\left(z^{\prime}\right)$ is continuous on $\gamma$, let

$$
F_{m}(z)=\int \frac{w\left(z^{\prime}\right)}{\left(z^{\prime}-z\right)^{m}} \frac{d z^{\prime}}{2 \pi i}
$$

for $m \geq 1$, then $F_{m}(z)$ is homolomorphic inside (resp. outside) $\gamma$ if $z$ is inside (resp. outside), and

$$
\frac{d F_{m}}{d z}=m F_{m+1}(z)
$$

Together we have given $f$ holomorphic then

$$
\frac{d^{m} f}{d z^{m}}=\frac{m!}{2 \pi i} \int \frac{f\left(z^{\prime}\right)}{\left(z^{\prime}-z\right)^{m+1}} d z
$$

and if $f$ is holomorphic on the ball of radius $R$ centred at $z_{0}$ then $f(z)=$ $\sum c_{k}\left(z-z_{0}\right)^{k}$ absolutely convergent in any ball $\left|z-z_{0}\right| \leq r \leq R$, a Taylor series with radius of convergence $R$. Because of this we say $f$ is analytic, meaning there exists a Taylor series with finite radius, and holomorphic is the same as analytic. The real counterpart is obviously false by considering the smooth function $f(x)=e^{-1 / x^{2}}, x \in \mathbb{R}$ which does not have a Taylor series with positive radius of convergence at the origin.

### 1.2 Convergence in the space of holomorphic functions

Let $H(B)$ be the space of functions $f: B \rightarrow \mathbb{C}$ analytic. Given a sequence of functions $\left\{f_{n}\right\}_{n \in \mathbb{N}}$ where $f_{n} \in H(B)$, what does it mean that $f_{n} \rightarrow f \in H(B)$ ? The correct notion is locally uniform convergence.

Recall that if $K \subseteq \mathbb{C}$ is compact then $\left(C(K),\|\cdot\|_{\infty}\right)$ is complete.

Lemma 1.1. Let $B \subseteq \mathbb{C}$ be a domain. There exists a sequence $\left\{K_{n}\right\}_{n \in \mathbb{N}}$ of compact sets contained in $B$ such that

1. $B=\bigcup_{n=0}^{\infty} K_{n}$,
2. $K_{n}$ is contained in the interior of $K_{n+1}$,
3. for all $K \subseteq B$ compact, exists $m$ such that $K \subseteq K_{m}$.

Lemma 1.2. Let $(X, d)$ be a metric space. Define

$$
\begin{aligned}
\mu: X \times X & \rightarrow[0,1) \\
(x, y) & \mapsto \frac{d(x, y)}{1+d(x, y)}
\end{aligned}
$$

then $\mu$ is a distance and $d, \mu$ define the same topology on $X$.
For $f, g \in C(B)$, define $d_{m}(f, g)=\|f-g\|_{K_{m}, \infty}$, then

$$
\rho(f, g)=\sum_{n=0}^{\infty} 2^{-n} \frac{d_{n}(f, g)}{1+d_{n}(f, g)}
$$

Theorem 1.3. $(C(B), \rho)$ is a complete metric space. Given $\left\{f_{m}\right\} \subseteq C(B)$, $f_{m} \rightarrow f$ in metric $\rho$ if and only if $f_{m} \rightarrow f$ uniformly on every compact subset of $B$, or equivalently on every $K_{n}$.

Theorem 1.4 (Weierstrass). $H(B) \subseteq C(B)$ is closed in $C(B)$. In particular it is complete.

### 1.3 Systems of ODEs

Let $D \subseteq \mathbb{C}^{n+1}$ be a domain. We denote a point in $D$ by $\left(z, y_{1}, \ldots, y_{n}\right)$ and write an analytic function $f_{j}: D \rightarrow \mathbb{C}$ as $f_{j}\left(z, y_{1}, \ldots, y_{n}\right)$ or simply $f_{j}(z, y)$.

Definition (system of ODEs). A system of ODEs of the first order and dimension $n$

$$
y^{\prime}=f(z, y)
$$

meaning

$$
\begin{aligned}
& y_{1}=f_{1}\left(z, y_{1}, \ldots, y_{n}\right) \\
& \vdots \\
& y_{n}=f_{n}\left(z, y_{1}, \ldots, y_{n}\right)
\end{aligned}
$$

is the problem of finding a domain $B \subseteq \mathbb{C}$ and $y: B \rightarrow \mathbb{C}^{n}$ analytic such that

1. $(z, y(z)) \in D$ and
2. $y^{\prime}(z)=f(z, y(z))$.

Definition (initial value problem). Let $\left(z_{0}, y_{0}\right) \in D$. An initial value problem (ivp) is the problem of finding $B \ni z_{0}$ and a solution $y(z)$ such that $y\left(z_{0}\right)=y_{0}$. It is usually written as

$$
y^{\prime}=f(z, y), y\left(z_{0}\right)=y_{0}
$$

Remark. Suppose $y=\left(y_{1}, \ldots, y_{n}\right)$. Then

$$
|y|=\sum_{j=1}^{n}\left|y_{j}\right|
$$

is a norm.

### 1.4 Existence and uniqueness theorem

Theorem 1.5 (existence and uniqueness theorem). Suppose $f: D \rightarrow \mathbb{C}^{n}$ is a vector valued function on a domain that is analytic and bounded, where

$$
D=\left\{(z, y) \in \mathbb{C}^{n+1}:\left|z-z_{0}\right|<a,\left|y-y_{0}\right|<b\right\}
$$

is a polydisc. Let $M=\sup _{z \in D}|f(z)|$. The the ivp

$$
y^{\prime}=f(z, y), y\left(z_{0}\right)=y_{0}
$$

has a unique solution $y(z)$ analytic on $U_{\alpha}\left(z_{0}\right)=\left\{z \in \mathbb{C}:\left|z-z_{0}\right|<\alpha\right\}$ where $\alpha=\min \{a, b / M\}$.

Exercise. Let $y(z)=y_{0}+\int_{z_{0}}^{z} f(\zeta, y(\zeta)) d \zeta$.

1. If $y(z)$ is an anlytic solution of the ivp then it also solves I.
2. Conversely if $y(z)$ is an analytic solution of I then it also solves the ivp.

Proof. The proof uses Picard's method of successive approximation: let

$$
\begin{aligned}
y_{0}(z) & =y_{0} \\
y_{m+1}(z) & =y_{0}+\int_{z_{0}}^{z} f\left(\zeta, y_{n}(\zeta)\right) d \zeta .
\end{aligned}
$$

We first show $y_{n}(z)$ is analytic on $U_{\alpha}\left(z_{0}\right) . y_{1}$ is analytic because the integrand is analytic (does not depend on the path) and

$$
\left|y_{1}(z)-y_{0}\right| \leq \int_{z_{0}}^{z}\left|f\left(\zeta, y_{0}\right)\right| d|\zeta| \leq M \cdot\left|z-z_{0}\right|<b
$$

Inductively $y_{n}$ is analytic and $\left|y_{n}(z)-y_{0}(z)\right| \leq M\left|z-z_{0}\right|<b$.
Next for every $K \subseteq D$ compact, the partial derivatives $\frac{\partial f_{i}(z, y)}{\partial y_{j}}$ are bounded on $K$. Then $f$ is Lipschitz continuous with respect to $y$, i.e. exists $k>0$ such that

$$
|f(z, y)-f(z, \tilde{y})| \leq k|y-\tilde{y}| .
$$

Let $K(\varepsilon)=\left\{\left|z-z_{0}\right| \leq \alpha-\varepsilon\right\}$. Then

$$
\left|y_{n}(z)-y_{0}\right| \leq M\left|z-z_{0}\right| \leq b-\varepsilon M .
$$

Now using telescopic sum

$$
y_{m}(z)=y_{0}+\sum_{j=0}^{m-1}\left(y_{j+1}(z)-y_{j}(z)\right)
$$

We can estimate the sum by

$$
\begin{aligned}
\left|y_{2}-y_{1}\right| & \leq \int_{z_{0}}^{z} \mid f\left(\zeta, y_{1}(\zeta)-f\left(\zeta, y_{0}\right)|d| \zeta \mid\right. \\
& \left.\leq k \int_{z_{0}}^{z} \mid y_{1}(\zeta)-y_{0}\right) d|\zeta| \\
& \leq k M \int_{z_{0}}^{z}\left|\zeta-z_{0}\right| d|\zeta| \\
& =\frac{M}{k} \frac{k^{2}\left|z-z_{0}\right|^{2}}{2}
\end{aligned}
$$

using Lipschitz continuity. By induction we can show the sum coverges. Thus on $K(\varepsilon), y_{m}(z)$ converges uniformly to some $y(z)$.

Next we show $y(z)$ on $K$ is an analytic solution of the ivp, or equivalently the corresponding integral equation

$$
y(z)=y_{0}+\int_{z_{0}}^{z} f(\zeta, y(\zeta)) d \zeta
$$

For reality check, note that $\left|y(z)-y_{0}\right|=\lim _{n \rightarrow \infty}\left|y_{n}(z)-y_{0}\right|<b$ so the expression $f(z, y(z))$ makes sense. By uniform convergence

$$
\lim _{n \rightarrow \infty} f\left(z, y_{n}(z)\right)=f(z, y(z))
$$

so

$$
\begin{aligned}
y(z) & =\lim _{n \rightarrow \infty} y_{n+1}(z) \\
& =y_{0}+\lim _{n \rightarrow \infty} \int_{z_{0}}^{z} f\left(\zeta, y_{n}(\zeta)\right) d \zeta \\
& =y_{0}+\int_{z_{0}}^{z} \lim _{n \rightarrow \infty} f\left(\zeta, y_{n}(\zeta)\right) d \zeta \text { uniform convergence } \\
& =y_{0}+\int_{z_{0}}^{z} f(\zeta, y(\zeta)) d \zeta
\end{aligned}
$$

For unqiueness, suppose $Y(z)$ is another solution defined on $\}\left|z-z_{0}\right| \leq$ $\left.\alpha-\varepsilon-\varepsilon^{\prime}\right\}$ such that $(z, Y(z)) \in D$. We show that on this domain $Y(z)=$ $\lim _{n \rightarrow \infty} y_{n}(z)$. This can be done using the integral equation to estimate $\mid Y(z)-$ $y_{n}(z) \mid$ inductively and show the limit is zero. Thus $y(z)$ and $Y(z)$ conincides on $\left|z-z_{0}\right| \leq \alpha-\varepsilon-\varepsilon^{\prime}$, so $y(z)$ is the analytic continuation of $Y(z)$.

To move from compact subsets to the entire domain, let $\left\{\varepsilon_{m}\right\}$ be a decreasing sequence converging to 0 and let $K_{m}=\left\{\left|z-z_{0}\right| \leq \alpha-\varepsilon_{m}\right\}$. Then $\left\{K_{m}\right\}$ is a candidate for the sets in lemma 1. By construction

$$
y_{n}^{K_{m}}(z)=\left.y_{n}^{K_{m+1}}(z)\right|_{K_{m}}
$$

so as $n \rightarrow \infty$,

$$
y^{K_{m}}(z)=\left.y^{K_{m+1}}\right|_{K_{m}}
$$

so $y^{K_{m+1}}$ is the analytic continuation of $y^{K_{m}}$ to $K_{m+1}$. Thus

$$
y_{n+z}(z)=y_{0}+\int_{z_{0}}^{z} f\left(\zeta, y_{n}(\zeta)\right) d \zeta
$$

converges on each $K_{m}$ uniformly, so by lemma 3 it converges to an analytic solution $y(z)$. Thus $y(z)$ is the unique solution on $U_{\alpha}\left(z_{0}\right)$.

## 2 Linear systems of ODEs

Consider a matrix-valued functions

$$
\begin{aligned}
A: U_{a}\left(z_{0}\right) & =\left\{z \in \mathbb{C}:\left|z-z_{0}\right|<a\right\} \rightarrow \operatorname{Mat}(n, \mathbb{C}) \\
b: U_{a}\left(z_{0}\right) & \rightarrow \mathbb{C}^{n}
\end{aligned}
$$

with each entry analytic. We are interested in a linear system of ODEs

$$
y^{\prime}(z)=\underbrace{A(z) y+b(z)}_{f(z, y)} .
$$

We use the matrix norm $|A|=\sum\left|A_{i j}\right|$ which can be easily checked to satisfy $|A y| \leq|A||y|$.

Lemma 2.1. Suppose $A, b: B \rightarrow \mathbb{C}$ be analytic and $K \subseteq B$ compact. Then $f(z, y)$ is bounded on $K$ and Lipschitz continuous with respect to $y$.

Proof.

$$
|f(z, y(z))-f(z, \tilde{y}(z))|=|A(z)(y-\tilde{y})| \leq|A(z)||y-\tilde{y}| \leq k|y-\tilde{y}|
$$

Theorem 2.2. Suppose $A, b$ are analytic on $U_{a}\left(z_{0}\right)$. Then the ivp

$$
y^{\prime}=A(z) y+b(z), y\left(z_{0}\right)=y_{0}
$$

where $y_{0}$ is any complex number, has a unique analytic solution in $U_{a}\left(z_{0}\right)$.
The proof is essentially the same, but note that there is no bound on $y_{0}$, and the domain of the solution has the same radius as the domain of $A$ and $b$. This is because $f$ is analytic on $D=\left\{z:\left|z-z_{0}\right|<\alpha\right\} \times \mathbb{C}$ so there is no bound on the second variable.

Theorem 2.3. Let $A, b$ be analytic on $B$, a simply connected domain. Then the ivp

$$
y^{\prime}=A(z) y+b(z), y\left(z_{0}\right)=y_{0}
$$

has unique solution $y(z)$ analytic on $B$.
Suppose $f_{1}: U_{1} \rightarrow \mathbb{C}$, $f_{2}: U_{2} \rightarrow \mathbb{C}$ are analytic, $U_{1} \cap U_{2} \neq \emptyset$ and $f_{1}(z)=f_{2}(z)$ for $z \in U_{1} \cap U_{2}$ then $f_{2}$ is the analytic continuation of $f_{1}$ on $U_{2}$.

Suppose $f: U \rightarrow \mathbb{C}$ is anlytic and $\gamma$ is a path from $z_{0} \in U$ to $\tilde{z}$. Then we say $f$ has an analytic continuation along $\gamma$ if we can cover $\gamma$ with a finite number of open balls such that $f$ has analytic continuation on the union of the balls.

Theorem 2.4 (monodromy theorem). Suppose $B$ is a connected domain. Suppose $f$ is analytic on an open ball $U \subseteq B$. Assume that $f$ has analytic continuation along any curve in $B$. Then if $\gamma_{1}$ and $\gamma_{2}$ are two homotopic path from $z_{0} \in U$ to $\tilde{z}$ then the analytic continuations $f_{\gamma_{1}}, f_{\gamma_{2}}$ agree on a neighbourhood of $\tilde{z}$.

In particular if in addition $B$ is simply connected then $f$ has a unique analytic continuation on the whole $B$.
Proof of Theorem 2.3. Suppose $\tilde{z} \in B$ and take a path $\gamma$ from $z_{0}$ to $\tilde{z}$. Note that as $B$ is open there exists $\varepsilon>0$ such that $B(z, \varepsilon) \subseteq B$ for all $z \in \operatorname{im} \gamma$. Cover the compact path $\gamma$ with finitely many open balls and for sufficiently small $\varepsilon^{\prime}$ we can solve the ivp on $B\left(z^{\prime}, \varepsilon^{\prime}\right)$. The solution is unique

### 2.1 Homogeneous system

A homogeneous system is one in which $b=0$, i.e. $y^{\prime}=A(z) y$. Let us asume that $A$ is analytic on $B$ simply connected. Then the system defines a linear operator

$$
\begin{aligned}
E: H(B) & \rightarrow H(B) \\
y & \mapsto y^{\prime}-A(z) y
\end{aligned}
$$

Then the space of solutions to the ODEs is ker $E$.
Theorem 2.5. If the system has dimension $n$ then $\operatorname{dim} \operatorname{ker} E=n$.
Proof. Let $v_{1}, \ldots, v_{n}$ be a basis of $\mathbb{C}^{n}$. Define $y_{i}(z)$ to be the unique analytic solution on $B$ of $y_{i}\left(z_{0}\right)=v_{i}$. ker $E$ is spanned by $y_{i}$ 's. Linear independence can be checked at $z_{0}$.
$y_{1}(z), \ldots, y_{n}(z)$ are called a fundamental system. We define $Y(z)$ to be the matrix whose columnes are $y_{i}(z)$. Then

$$
\frac{d Y}{d z}=\left[\frac{d y_{1}}{d z}|\cdots| \frac{d y_{n}}{d z}\right]=\left[A y_{1}|\cdots| A y_{n}\right]=A\left[y_{1}|\cdots| y_{n}\right]=A(z) Y(z)
$$

$Y(z)$ is called a fundamental matrix solution.
Corollary 2.6. Analytic continuation along any curve preservse linear independence.

In other words, $\operatorname{det} Y\left(z_{0}\right) \neq 0$ implies $\operatorname{det} Y(z) \neq 0$ for all $z \in B$. Suppose we have another fundamental matrix solution

$$
\tilde{Y}(z)=\left[\tilde{y}_{1}(z)|\cdots| \tilde{y}_{n}(z)\right]=\left[\sum_{k=1}^{n} c_{k 1} y_{k}(z)|\cdots| \sum_{k=1}^{n} c_{k n} y_{k}(z)\right]=Y(z) C
$$

where $C=\left(c_{i j}\right)$ nonsingular.
One application: suppose we have a homogeneous ODE of order $n$

$$
u^{(n)}+p_{1}(z) u^{(n-1)}+\cdots+p_{n}(z) u=0
$$

where $u$ is a scalar and $p_{i}: B \rightarrow \mathbb{C}$ analytic. Then this can be recast into a linear ODE by setting $y_{1}=u, y_{2}=y^{\prime}$ etc:

$$
y^{\prime}=\left(\begin{array}{cccc}
0 & 1 & & \\
& 0 & 1 & \\
\vdots & & & \\
-p_{n} & -p_{n-1} & \cdots & -p_{1}
\end{array}\right) y
$$

This is called the companion matrix. $u_{1}(z), \ldots, u_{n}(z)$ is a fundamental system.

### 2.2 Inhomogeneous system

We will see that it suffices to study homogeneous case when studying linear systems. Suppose we have an inhomogeneous problem $y^{\prime}=A(z) y+b(z)$, where $A, b$ are analytic on a simply connected domain $B$. Then in matrix form it is

$$
\frac{d W}{d z}=A(z) W+[b(z)|\cdots| b(z)]
$$

where $W$ is an (unknown) $n \times n$ matrix. Of course it is a special case of

$$
\begin{equation*}
\frac{d W}{d z}=A(z) W+[b(z)|\cdots| b(z)] \tag{NH}
\end{equation*}
$$

where $F$ is analytic and matrix valued. We call $\frac{d Y}{d z}=A(z) Y$ the associated homogeneous system, and let its fundamental matrix system be $Y(z)$ with $\operatorname{det} Y(z) \neq 0$.

Theorem 2.7. All solutions of (NH) are

$$
W(z)=Y(z)\left[C+\int_{z_{0}}^{z} Y^{-1}(\zeta) F(S) d \zeta\right]
$$

where $Y(z)$ is a fundamental solution of the homogeneous system.
The ivp $W\left(z_{0}\right)=W_{0}$ has solution obtained with $C=Y\left(z_{0}\right)^{-1} W_{0}$.
Proof. Differentiate the purported solution,

$$
\begin{aligned}
W^{\prime} & =Y^{\prime}\left[C+\int_{z_{0}}^{z} Y^{-1}(\zeta) F(\zeta) d \zeta\right]+Y(z) Y^{-1}(z) F(z) \\
& =A(z) Y(z)\left[C+\int_{z_{0}}^{z} Y^{-1}(\zeta) F(\zeta) d \zeta\right]+F(z) \\
& =A(z) W(z)+F(z)
\end{aligned}
$$

Every solution of (NH) is obtained from some ivp $W_{0}$, from which we derive $C=Y\left(z_{0}\right)^{-1} W_{0}$. Therefore we obtain every solution.

The formula is obtained by "variation of parameter". In fact it is instructive to do the following easy exercise: we seek a solution of the form $W(z)=Y(z) C(z)$. Substitute into (NH), we will get $C^{\prime}(z)=Y^{-1}(z) F(z)$. Integrate to get the desired solution in the theorem.

Now back to the homogeneous problem $\frac{d Y}{d z}=A(z) Y$. Suppose $A$ is holomorphic on $B$. Take $z_{0} \in B, r>0$ maximal such that $U_{r}\left(z_{0}\right) \subseteq B$. In particular $A$ is holomorphic in the ball $U_{r}\left(z_{0}\right)$. $A$ has a power series representation $A(z)=\sum A_{j}\left(z-z_{o}\right)^{j}$ convergent uniformly in every compact subset of $U_{r}\left(z_{0}\right)$ and the radius of convergence is $r$. What we know is that

1. there exists a unique solution to the ivp $Y\left(z_{0}\right)=Y_{0}$.
2. $Y(z)$ is analytic in $U_{r}(a)$.

Thus $Y(z)=\sum Y_{k}\left(z-z_{0}\right)^{k}$ for $\left|z-z_{0}\right|<r$. Substitute into the equation,

$$
\sum_{k=1}^{\infty} k Y_{k}\left(z-z_{0}\right)^{k-1}=\sum_{j=0}^{\infty} \sum_{k=0} A_{j} Y_{k}\left(z-z_{0}\right)^{j+k}
$$

i.e.

$$
\sum_{m=0}^{\infty}(m+1) Y_{m+1}\left(z-z_{0}^{m}=\sum_{m=0}^{\infty}\left(\sum_{k=0}^{m} A_{m-k} Y_{k}\right)\left(z-z_{0}\right)^{m}\right.
$$

We get the recurrence relation

$$
(m+1) Y_{m}=\sum_{k=0}^{m} A_{m-k} Y_{k}
$$

with (arbitrary) initial condition $Y_{0}$. Up to order $M$, the solution has local representation

$$
Y(z)=\sum_{k=0}^{M} Y_{k}\left(z-z_{0}^{k}+O\left(\left(z-z_{0}\right)^{M+1}\right)\right.
$$

## 3 Singularities and monodromies

### 3.1 Classification of isolated singularities

Suppose $f: B \backslash\{a\} \rightarrow \mathbb{C}$ is holomorphic where $B$ is a domain. Then $a$ is called an isolated singularity. By basic complex analysis $a$ is one of the below

1. $a$ is a removable singularity if $|f(z)|$ is bounded in some $U_{r}(a) \backslash\{a\}$, in which case the limit $\lim _{z \rightarrow a} f(z)$ exists, and $\tilde{f}(z)= \begin{cases}f(z) & z \neq a \\ \lim _{z \rightarrow a} f(z) & z=a\end{cases}$ is holomorphic on $B$.
2. $a$ is a pole if $\lim _{z \rightarrow \infty}|f(z)|=\infty$, in which case exists $r$ such that $f(z)=$ $\frac{g(z)}{(z-a)^{m}}$, where $g$ is holomorphic on $U_{r}(a)$ and $g(a) \neq 0$ and $m \geq 1 . m$ is called the order of the pole. $f(z)$ has a Laurent exptension at $z=a$.
3. $a$ is an essential singularity if $\lim _{z \rightarrow a}|f(z)|$ does not exist, in which case $f(z)$ takes all possible complex values except possibly one (called the lacunary value) in a neighbourhood of $a$.

### 3.2 Singularities of linear systems

Now consider $y^{\prime}=A(z) y+b(z)$ where $A, b$ are analytic on $B \backslash\left\{a_{0}, \ldots, a_{m}\right\}$. Let $\overline{\mathbb{C}}=\mathbb{C} \cup\{\infty\}$. Then for a function $f(z)$ defined on $\mathbb{C}$, we can consider $\infty$ as an isolated singularity and use the classification before.

We call $a_{1}, \ldots$ singularities of the linear differential system.

Theorem 3.1. Given a solution $y(z)$ of a linear system, the set of singularities of $y$ is a subset of the singularities of the equation.

Proof. If $A, b$ are analytic in a ball (which is simply connected) then all $y(z)$ are analytic there too. If $z=a$ is a singularity of $y(z)$, then necessarily $z=a$ is a singularity of either $A(z)$ or $b(z)$.

## Remark.

1. The converse to the theorem is false. Consider $y^{\prime}=\frac{\mu}{z-a} y$ where $\mu \in \mathbb{C}$. The solution $y=c(z-a)^{\mu}$ is
(a) regular at $z=a$ if $\mu \in \mathbb{N}$,
(b) a pole if $\mu \in\{-1,-2, \ldots\}$,
(c) a branch point if $\mu \in \mathbb{C} \backslash \mathbb{Z}$.
2. Note that if we consider $A, b$ as functions defined on $\overline{\mathbb{C}}$, then $z=\infty$ can be a singularity of the solution even if it is not a singularity of the equation (i.e. $A, b$ have a removable singularity at $\infty$ ). For example consider $y^{\prime}=y$ : solution $y(z)=c e^{z} . z=\infty$ is a singularity.
3. The theorem only applies to linear ODEs. Consider for example $y^{\prime}=-y^{2}$. Solution $y(z)=\frac{1}{z-a} \cdot z=a$ is a pole of $y(z)$ but not a singularity of the equation.

### 3.3 Monodromy for homogeneous systems

Consider the homogeneous problem $\frac{d Y}{d z}=A(z) Y$ where $A: B \backslash\{a\} \rightarrow \mathbb{C}$ analytic. We have seen that analytic continuation along $\ell$ from $z_{0}$ to $z^{\prime}$ defines $y(z)$ analytic in a neighbourhood of $z^{\prime}$. Take another path $\tilde{\ell}$ which is not homotopic to $\ell$ and analytic continuation along $\ell^{\prime}$ defines $\tilde{y}(z)$.

Note that if $B$ is a disc, we can take a branch cut $L$ from $a$ to the boundary of $B$, and then $B \backslash L$ is simply connected. $y, \tilde{y}$ are analytic on $B \backslash L . y, \tilde{y}$ are called branches of each other. Let $\gamma=\ell^{-1} \cdot \tilde{\ell}$, then travelling along $\gamma$ transforms $y$ to $\tilde{y}$. $\tilde{y}$ is an analytic continuation of $y$.

If there are more than one singularity $\left\{a_{1}, \ldots, a_{m}\right\}$ in $B$, we can do a branch cut $L_{i}$ from each $a_{i}$, and $y, \tilde{y}$ are analytic on $B \backslash\left\{L_{1}, \ldots, L_{m}\right\}$.

For a homogeneous system $\frac{d Y}{d z}=A(z) Y$, we proved that there exsits a fundamental solution $Y(z)=\left[y_{1}(z)|\cdots| y_{n}(z)\right]$ analytic on $\mathbb{C} \backslash\left\{L_{1}, \ldots, L_{n}\right\}$. If we take another path we get $\tilde{Y}(z)$. But columns of $Y$ form a basis, we can express

$$
\tilde{Y}(z)=\left[\sum m_{i 1} y_{i}|\cdots| \sum_{i} m_{n i} y_{i}\right]=Y(z) M_{\gamma}
$$

where $M_{\gamma}=\left(m_{i j}\right)$ is called the monodromy matrix of $Y(z)$ associated to (the homotopy class) of $\gamma$.

Recall from algebraic topology that $\pi_{1}\left(\overline{\mathbb{C}}\left\{a_{1}, \ldots, a_{m}, \infty\right\}, z_{0}\right)$ is the free group generated by $\gamma_{1}, \ldots, \gamma_{m}$. Also there is an automorphism given by conjugations

$$
\begin{aligned}
\pi\left(\overline{\mathbb{C}} \backslash\left\{a_{1}, \ldots, a_{m}, \infty\right\}, z_{0}\right) & \rightarrow \pi\left(\overline{\mathbb{C}} \backslash\left\{a_{1}, \ldots, a_{m}, \infty\right\}, z_{1}\right) \\
\gamma & \mapsto \lambda \gamma \lambda^{-1}
\end{aligned}
$$

where $\lambda$ is a path from $z_{1}$ to $z_{0}$. Now suppose $Y\left(z_{1}\right)$ is the analytic continuation of $Y\left(z_{0}\right)$ along $\lambda^{-1}$. Then obviously the analytic continuations of $Y\left(z_{0}\right)$ along $\gamma \cdot \lambda^{-1}$ and along $\lambda^{-1} \cdot\left(\lambda \cdot \gamma \cdot \gamma^{-1}\right)$ yields the same result

$$
\begin{array}{r}
Y\left(z_{0}\right) \xrightarrow{\gamma} Y\left(z_{0}\right) M_{\gamma} \\
\stackrel{\lambda^{-1}}{\downarrow} \xrightarrow{\downarrow \lambda^{-1}} \\
Y\left(z_{1}\right) \xrightarrow{\lambda \gamma \lambda^{-1}} Y\left(z_{1}\right) M_{\gamma}
\end{array}
$$

But the bottom row is precisely the monodromy transformation associated to $\lambda \gamma \lambda^{-1}$ so we have $Y(z) M_{\gamma}=Y(z) M_{\lambda \gamma \lambda^{-1}}$. Since $Y(z)$ is invertible, $M_{\lambda \gamma \lambda^{-1}}=$ $M_{\gamma}$.

Theorem 3.2. Let $z_{0}, Y_{0}$ be given ( $\operatorname{det} Y_{0} \neq 0$ ) and let $Y(z)$ be the unique fundamental matrix solution analytic in $\mathbb{C} \backslash\left\{L_{1}, \ldots, L_{m}\right\}$ solving the ivp.

$$
\begin{aligned}
\phi: \pi_{1}\left(\overline{\mathbb{C}} \backslash\left\{a_{1}, \ldots, a_{m}, \infty\right), z_{0}\right) & \rightarrow \operatorname{GL}(n, \mathbb{C}) \\
\gamma & \mapsto M_{\gamma}
\end{aligned}
$$

this is called the monodromy representation of $\pi_{1}$. This is an anti-representation.
Proof. Exercise.
$\gamma_{1}, \ldots, \gamma_{m}$ generates the group $\mathcal{M}=\left\langle M_{\gamma_{1}}, \ldots, M_{\gamma_{n}}\right\rangle$, the monodromy group associated with the ivp $Y\left(z_{0}\right)=Y_{0}$.

If we have another ivp $Y\left(z_{0}\right)=\hat{Y}_{0}$, ther is a unique solution $\hat{Y}(z)$ defined on $\mathbb{C} \backslash\left\{L_{1} \cup \cdots L_{m}\right\}$. The two solutions are related by $\hat{Y}(z)=Y(z) C$ where $C=Y_{0}^{-1} \hat{Y}_{0} \in \mathrm{GL}(n, \mathbb{C})$. Thus along the path $\gamma$,

$$
\hat{Y}(z)=Y(z) C \mapsto Y(z) M_{\gamma} \cdot C=Y(z) C \cdot \underbrace{C^{-1} M_{\gamma} C}_{\hat{M}_{\gamma}} .
$$

Thus from the data of $A(z)$, we get $\left\langle C^{-1} M_{\gamma_{1}} C, \ldots, C^{-1} M_{\gamma_{n}} C\right\rangle$ for all $C \in$ $\mathrm{GL}(n, \mathbb{C})$. This is an equivalence relation on monodromy groups.

Definition (direct monodromy map). Fix points $a_{1}, \ldots, a_{m} \in \mathbb{C}$. The direct monodromy map is the map
$\left\{A(z)\right.$ isolated singularities $\left.a_{1}, \ldots, a_{m}, \infty\right\} \rightarrow\{$ monodromy group $\} /$ conjugation
It is worth remarking here the opposite problem, that is, given monodromy groups up to conjugation, can we determine the differential system? This is called the inverse monodromy problem or Riemann-Hilbert problem.

### 3.4 Solutions on universal covers

Fix a base point $z_{0}$. Analytic continuation gives a function $Y(z)=" y(z, \ell) "$ that depends on (the homotopy class of) the path $\ell$. Define an equivalence relation $(z, \ell) \sim\left(z^{\prime}, \ell^{\prime}\right)$ if and only if $z=z^{\prime}, \ell \simeq \ell^{\prime}$. The the set of equivalence classes is in bijection with the universal cover. The fundamental group acts on the universal cover by deck transformation. If there is no chance of confusion we denote a point $[(z, \ell)]$ by $\tilde{z}$.

The analytic continuation of $Y(z)$ along any path $\ell$ gives a function $\mathbb{Y}$ on the universal convering so, defines a function $\mathbb{Y}(\tilde{z})$ is a function of $\tilde{z}$ on the universal cover. $Y(z)$ is a branch of $\mathbb{Y}(\tilde{z})$.

We have seen that $\pi_{1}\left(\overline{\mathbb{C}} \backslash\{0, \infty\} ; z_{0}\right) \cong \mathbb{Z}$. A point $[(z, \ell)] \in \overline{\mathbb{C}} \backslash\{0, \infty\}$ can be represented by $\left(z, \ell_{m}=\gamma^{m} \cdot \ell\right)$ where $m \in \mathbb{Z}$ and $\ell$ is a path from $z_{0}$ to $z$. We give the another presentation. Take a branch cut $(-\infty, 0]$ in the complex domain and consider the principal branch of logarithm, defined so that $\ln z \in \mathbb{R}$ for $z>0$. Then $z \mapsto \ln z$ represents this: $\ln z=\ln |z|+i \arg z$. If a path crosses the branch cut $m$ times, the $\arg z$ is transformed to $\arg z+2 \pi m$. In this case the universal covering is the exponential map

$$
\begin{aligned}
\exp : \overline{\mathbb{C}} \backslash\{0, \infty\} & \rightarrow \overline{\mathbb{C}} \backslash\{0, \infty\} \\
& \ln z+2 \pi i m
\end{aligned}>z
$$

For this reason $\overline{\mathbb{C}} \backslash\{0, \infty\}$ is called the Riemann surface of the logarithm. By a slightly confusing change of notation, we can represent a point in the universal cover by $\tilde{z}=|z| e^{i \arg z}$, where the exponential is a formal symbol (so $|z| e^{i \arg z}$ and $|z| e^{i \arg z+2 \pi i m}$ are different points). The advantage of this notation is that monodromy can be computed easily just by symbolic manipulation. For example

$$
\tilde{z}^{c}=e^{c \ln \tilde{z}}=|z|^{c} e^{i c \arg z}
$$

whose meaning is that if the locally defined function $z \mapsto z^{c}$ transforms with a multiple of $e^{2 \pi i c}$ when taken along the path that loops around the origin once.

## 4 Classification of isolated singularities

### 4.1 Matrix exponential and logarithm

We endow $\operatorname{Mat}(n, \mathbb{C})$ with the matrix norm $|A|=\sum\left|A_{i j}\right|$. It is a complete metric space. Then

$$
e^{A}=\sum_{k=0}^{\infty} \frac{A^{k}}{k!}
$$

is convergent. Similarly for $f$ analytic, we can define

$$
e^{A \cdot f(z)}=\sum_{k=0}^{\infty} \frac{A^{k} f(z)^{k}}{k!}
$$

which is locally uniformly convergent so analtyic. As a special case, take the principal branch of logarithm and we can define

$$
z^{A}=e^{A \log z}
$$

Example. If going along a path transforms $z \mapsto z \cdot e^{2 \pi i}$ then

$$
z^{A}=e^{A \ln z} \mapsto e^{2 \pi i A} z^{A}
$$

and indeed $e^{2 \pi i A}$ is the monodromy matrix. We can think of this as defined on the universal covers by $\tilde{z}^{A}=|z|^{A} e^{A i \arg z}$, where the formal exponential helps us keep track of monodromy actions.
Exercise. Matrix exponential has the following properties:

1. if $A=G B G^{-1}$ for some $G \in \mathrm{GL}(n, \mathbb{C})$ then $e^{A}=G e^{B} G^{-1}$. It follows that $z^{A}=G z^{B} G^{-1}$.
2. If $[A, B]=0$ then $e^{A+B}=e^{A} e^{B}$.
3. $\operatorname{det} e^{A}=e^{\operatorname{tr} A}$. In particular $\operatorname{det} e^{A} \neq 0$ (hint: put into Jordan form).
4. $\frac{d}{d z} e^{A f(z)}=A f^{\prime}(z) e^{A f(z)}$ so $\frac{d}{d z} z^{A}=\frac{A}{z} \cdot z^{A}$. Thus $\frac{d Y}{d z}=\frac{A}{z} Y$ has a fundamental solution $Y(z)=z^{A}$. As usual we write $\mathbb{Y}(\tilde{z})=\tilde{z}^{A}$ for the function defined on the universal cover.
For $k \geq 2, \frac{d Y}{d z}=\frac{A}{z^{k}} Y$ has solution $Y(z)=\exp \left(-\frac{A}{k-1} \frac{1}{z^{k-1}}\right)$. On the other hand $\frac{d Y}{d z}=z^{k} A Y$ for $k \geq 0$ has solution $Y(z)=\exp \left(\frac{A}{k+1} z^{k+1}\right)$.
5. If $\left[A_{i}, A_{j}\right]=0$ then

$$
\frac{d Y}{d z}=\left(\frac{A_{1}}{z-a_{1}}+\cdots+\frac{A_{m}}{z-a_{m}}\right) Y
$$

has a solution just as if they are numbers: $Y(z)=\left(z-a_{1}\right)^{A_{1}} \cdots\left(z-a_{m}\right)^{A_{m}}$.
6. Suppose all matrices commute and $h \geq 0, k \geq 2$, then

$$
\frac{d Y}{d z}=\left(\frac{1}{z^{k}}\left(A_{0}+A_{1} z+\cdots+A_{k-2} z^{k-2}\right)+\frac{B}{z}+z^{h}\left(C_{0}+\frac{C_{1}}{z}+\cdots+\frac{C_{h}}{z^{h}}\right)\right) Y
$$

has solution

$$
\mathbb{Y}(\tilde{z})=\tilde{z}^{B} \exp \left(\frac{A_{0}}{1-k} \frac{1}{z^{k-1}}+\cdots+\frac{A_{k-2}}{-z}\right) \exp \left(\frac{C_{0}}{h+1} z^{h+1}+\cdots+C_{h} z\right)=\Phi(z) \tilde{z}^{B}
$$

Definition. A logarithm of $A$, where $\operatorname{det} A \neq 0$, written $L=\ln A$, is a matrix such that $e^{L}=A$.

Matrix logarithm is not unique: suppose there exists an $L$, then $L+2 \pi i k I$ is also a logarithm.

We can compute matrix logarithm by making a series of reductions. Assuming everything exists, note

$$
e^{\ln A}=A=G J G^{-1}=G e^{\ln J} G^{-1}=e^{G \ln J G^{-1}}
$$

so we can let $\ln A=G \ln J G^{-1}$ (up to some issue of determinancy). Thus it is left to define and compute $\ln J$.
Exercise. Check that exp preserves block-diagonal form. Conversely suppose $e^{A}$ is block-diagonal then $A$ is also block diagonal.

Write $A=G J G^{-1}$ where $J$ is the Jordan normal form with Jordan blocks $J_{1}, \ldots, J_{s}$. Write $J_{i}=\lambda_{i} I_{i}+H_{i}$ where $\lambda_{i}$ is the eigenvalue of $J_{i}$ and $H=$ $\left(\begin{array}{ccc}0 & 1 & \\ 0 & 1 \\ & & \ddots\end{array}\right)$. The size of Jordan block $J_{i}$ is $m_{i}$. We define

$$
\begin{aligned}
\ln J_{i} & =\ln \left(\lambda_{i} I_{i}\left(I_{i}+\frac{H_{i}}{\lambda_{i}}\right)\right) \\
& =\ln \left(\lambda_{i} I_{i}\right)+\ln \left(I_{i}+\frac{H_{i}}{\lambda_{i}}\right) \\
& =\left(\ln \lambda_{i}\right) I_{i}+\sum_{k=0}^{\infty} \frac{(-1)^{k+1}}{\lambda_{i}}\left(\frac{H_{i}}{\lambda_{i}}\right)^{k} \\
& =\left(\ln \lambda_{i}\right) I_{i}+\sum_{k=0}^{m_{i}-1} \frac{(-1)^{k+1}}{\lambda_{i}}\left(\frac{H_{i}}{\lambda_{i}}\right)^{k}
\end{aligned}
$$

as $H_{i}$ is nilpotent.
Again note that $\ln J_{i}$ is defined up to integer multiples of $2 \pi i I_{i}$.

### 4.2 Local structure of fundamental matrices at isolated singularities

Consider $\frac{d Y}{d z}=A(z) Y$ where $A(z)$ is analytic on $U(a) \backslash\{a\}$. The universal covering can be represented by logarithm

$$
\ln (z-a)=\ln |z-a|+i \arg (z-a) .
$$

We do the following steps:

1. take a branch cut, define a fundamental solution $Y(z)$.
2. Take a loop so $(\tilde{z}-a) \mapsto(\tilde{z}-a) e^{2 \pi i}$, and the branch $Y(z) \mapsto Y(z) M$.
3. $\mathbb{Y}(\tilde{z})=Y(z) M^{k}$ (i.e. defined on all branches). Note that $M$ is also the action of an element of the fundamental group on $\mathbb{Y}$ (in deck transformation).

Theorem 4.1. Given the system $\frac{d Y}{d z}=A(z) Y$ where $A(z)$ is analytic on $U(a) \backslash\{a\}$, for all $\tilde{z} \in U \widetilde{(a) \backslash\{a\}}$, we have the following representation

$$
\mathbb{Y}(\tilde{z})=\Phi(z) \cdot(\tilde{z}-a)^{L}
$$

where $L=\frac{1}{2 \pi i} \ln M . \Phi(z)$ is single valued on $U(a) \backslash\{a\}$, $\operatorname{det} \Phi(z) \neq 0$ for $z \neq a$.

Proof. Define $\Phi(\tilde{z})=\mathbb{Y}(\tilde{z})(\tilde{z}-a)^{-L}$. It is single valued as when taken along a path,

$$
\Phi(\tilde{z}) \mapsto \mathbb{Y}(\tilde{z}) M \cdot \underbrace{e^{-2 \pi i L}}_{M^{-1}}(\tilde{z}-a)^{-L}=\mathbb{Y}(\tilde{z})(\tilde{z}-a)^{-L}=\Phi(\tilde{z}) .
$$

Also $\operatorname{det} \Phi(\tilde{z})=\operatorname{det} \mathbb{Y}(\tilde{z}) \operatorname{det}(\tilde{z}-a)^{-L} \neq 0$.
Example. The last exercise of the previous section is put in this form.
Now let's do some computation. Write $L=G J G^{-1}$ and so

$$
\mathbb{Y}(\tilde{z})=\Phi(z)(\tilde{z}-a)^{L}=\underbrace{\Phi(z) G}_{\Psi(z)}(\tilde{z}-a)^{J} G^{-1}
$$

Then $\tilde{\mathbb{Y}}(\tilde{z})=G \mathbb{Y}(\tilde{z})$ is another fundamental solution with structre

$$
\tilde{\mathbb{Y}}(\tilde{z})=\Psi(z)(\tilde{z}-a)^{J} .
$$

For simplicity suppose $a=0$, then

$$
\tilde{z}^{J_{i}}=z_{i}^{\lambda}\left(\begin{array}{cccccc}
1 & \ln z & \frac{(\ln z)^{2}}{2!} & \ldots & \frac{(\ln z)^{m_{i}-1}}{\left(m_{i}-1\right)!} & \\
& 1 & \ln z & & & \\
\vdots & & & & & \\
& & \ddots & & & 1
\end{array}\right)
$$

### 4.3 Regular singularities

Definition (regular singularity). Given $f: \widetilde{U_{r}} \widetilde{(a) \backslash}\{a\} \rightarrow \mathbb{C}$ analytic, $z=a$ is a regular singularity if exists $m \in \mathbb{R}_{+}$such that

$$
\lim _{\alpha<\arg (\underset{z}{z}-a)<\beta}(\tilde{z}-a)^{m} f(z)=0
$$

for all $\beta-\alpha<2 \pi$. Otherwise $z=a$ is an irregular singularity.
Easy to see that if $f, g$ has regular singularity at $z=a$ then so do $f+g$ and fg.

## Example.

1. $f(z)=\frac{\psi(\tilde{z})}{(\tilde{z}-a)^{r}}$ where $\psi(z)$ is analytic at $z=a, r \in \mathbb{R}$ has a regular singularity.
2. $f(z)=(\tilde{z}-a)^{\mu} \ln (\tilde{z}-a)^{\nu}$ where $\mu, \nu \in \mathbb{C}$ has a regular singularity.
3. $f(z)=\exp \left(\frac{1}{(\tilde{z}-a)^{\mu}}\right)$ where $\mu \in \mathbb{C}$ has an irregular singularity.
4. The solution $\mathbb{Y}(z)=\Phi(z)(\tilde{z}-a)^{L}$, (suppose $(\tilde{z}-a)^{L}$ is regular) depends on $\Phi(z)$. For example if $z=a$ a removable singularity or a pole of $\Phi(z)$ then $z=a$ is regular for $\mathbb{Y}(\tilde{z})$.

Definition. $z=a$ is a regular singularity for $\frac{d Y}{d z}=A(z) Y$ if exists a fundamental solution $\mathbb{Y}(\tilde{z})$ such that $z=a$ is regular for $\mathbb{Y}(\tilde{z})$.

Remark. Every other fundamental solution can be given by $\mathbb{Y}(\tilde{z}) C$ for $\operatorname{det} C \neq$ 0 so this is well-defined.

Theorem 4.2. Suppose $z=a$ is regular for $\frac{d Y}{d z}=A(z) Y$. Then in $\mathbb{Y}(\tilde{z})=$ $\Phi(z)(\tilde{z}-a)^{L}, z=a$ is a removable singularity or a pole for $\Phi(z)$, or a pole for $A(z)$.

Proof. $\Phi(z)=\mathbb{Y}(\tilde{z})(\tilde{z}-a)^{-L}$ so $z=a$ is regular for $\Phi(z)$. Unpack the definition, we can find $m \in \mathbb{N}$ such that

$$
\lim _{z \rightarrow a}(z-a)^{m} \Phi(z)=0
$$

for $\alpha<\tilde{z}-a<\beta$. But both $(z-a)^{m}$ and $\Phi(z)$ are single valued so we can cover the entire disc by sectors and remove the condition $\alpha<\tilde{z}-a<\beta$. Thus $\Phi(z)=\frac{1}{(z-a)^{m}} \cdot(z-a)^{m} \Phi(z)$ has at most a pole at $z=a$.

For the other condition,

$$
\begin{aligned}
A(z) & =\frac{d \mathbb{Y}(\tilde{z})}{d z} \mathbb{Y}(\tilde{z})^{-1} \\
& =\frac{d}{d z}\left(\Phi(z)(\tilde{z}-a)^{L}\right) \cdot(\tilde{z}-a)^{-L} \Phi(z)^{-1} \\
& =\frac{d \Phi(z)}{d z} \Phi^{-1}(z)+\Phi(z) \frac{L}{z-a} \Phi(z)^{-1}
\end{aligned}
$$

so $z=a$ is a pole.
Example. Consider

$$
A(z)=\frac{1}{z^{2}}\left(\begin{array}{ll}
0 & 1 \\
0 & 0
\end{array}\right)+\frac{1}{z}\left(\begin{array}{ll}
0 & 0 \\
0 & 1
\end{array}\right)
$$

Write the solution as $Y(z)=\left(y_{1}, y_{2}\right)^{T}$, then we have

$$
y_{1}^{\prime}=\frac{1}{z^{2}} y_{2}, y_{2}^{\prime}=\frac{1}{z} y_{2} .
$$

Solve to get $y_{1}(z)=C_{2} \ln z+C_{1}, y_{2}(z)=C_{2} z$. We choose $\left(c_{1}, c_{2}\right)=(1,0),(0,1)$ to get the fundamental solution

$$
\mathbb{Y}(\tilde{z})=\left(\begin{array}{cc}
1 & \ln \tilde{z} \\
0 & z
\end{array}\right)
$$

The monodromy associated to $z \mapsto z e^{2 \pi i}$ is

$$
\mathbb{Y}(\tilde{z}) \mapsto\left(\begin{array}{cc}
1 & \ln \tilde{z}+2 \pi i \\
0 & z
\end{array}\right)=\mathbb{Y}(\tilde{z}) \cdot e^{\left(\begin{array}{cc}
0 & 2 \pi i \\
0 & 0
\end{array}\right)}=\Phi(z) z\left(\begin{array}{ll}
0 & 1 \\
0 & 0
\end{array}\right)
$$

where $\Phi=\left(\begin{array}{cc}1 & 0 \\ 0 & z\end{array}\right)$ and $\operatorname{det} \Phi(z) \neq 0$ for $z \neq 0$.

### 4.4 Singularities of the first and second kind

Given the system $\frac{d Y}{d z}=A(z) Y$ where $A: U_{r}(a) \backslash\{a\} \rightarrow \operatorname{Mat}(n, \mathbb{C})$ analytic and $z=a$ is an isolated singularity of $A$ that is not an essential singularity. This means we can write $A(z)=\frac{\tilde{A}(z)}{(z-a)^{r+1}}$ where $\tilde{A}$ is analytic on $U_{R}(a)$ and $\tilde{A}(z) \neq 0$.

Definition (Fuchsian singularity). If $r=0$ then $z=a$ is called a singularity of the first kind (name used by Coddington-Levinson/Balsen-Jurkat-Lutz) or Fuchsian singularity. If $r \geq 1$ then $z=a$ is called a singularity of the second kind, or occassionally (and confusingly) called a irregular singularity. If $r \leq-1$ then it is not a singularity.

Theorem 4.3. If $z=a$ is a Fuchsian singularity then $z=a$ is a regular singularity.

Note that the converse is not true, as for example in the example above there is a pole of second order.

We can similarly classify singularities at infinity. Suppose $A(z)$ is analytic for $|z|>R$. Make the change of variable $z=\frac{1}{t}$ so $\frac{d}{d z}=-t^{2} \frac{d}{d t}$, we have

$$
\frac{d Y}{d t}=-\frac{A(1 / t)}{t^{2}} Y
$$

Thus $A_{*}(t)=A\left(\frac{1}{t}\right)$ is analytic for $0<|t|<\frac{1}{R}$ and $t=0$ is an isolated singularity. Write $-\frac{A_{*}(t)}{t^{2}}=\frac{\tilde{A}(t)}{t^{r+1}}$ so

$$
A_{*}(t)=-t^{-r+1} \tilde{A}(t)=-z^{r-1} \tilde{A}^{\prime}(z)
$$

so for $r=0$ we call it a Fuchsian singularity and similarly for $r \geq 1$ and $r \leq-1$.

### 4.5 Linear systems with rational coefficients

Now suppose $A(z)$ is defined on $\overline{\mathbb{C}} \backslash\left\{a_{1}, \ldots, a_{m}, \infty\right\}$ and suppose $a_{1}, \ldots, a_{m}, \infty$ are at most poles. The locally around $z=a_{j}$ we can write $A(z)=A^{(j)}(z)+$ $\operatorname{reg}\left(z-a_{j}\right)$ where the regular part is analytic at $z=a_{j}$ and $\operatorname{reg}\left(z-a_{j}\right)=O(1)$ as $z \rightarrow a_{j}$. In detail

$$
A^{(j)}(z)= \begin{cases}\frac{1}{\left(z-a_{j}\right)^{r_{j}+1}}\left(A_{0}^{(j)}+A_{1}^{(j)}\left(z-a_{j}\right)+\cdots+A_{r_{j}}^{(j)}\left(z-a_{j}\right)^{r_{j}}\right) & r_{j} \geq 0 \\ 0 & r_{j} \leq-1\end{cases}
$$

and for $z=\infty$ write $A(z)=A^{(\infty)}(z)+\operatorname{reg}\left(\frac{1}{z}\right)$ where $\operatorname{reg}\left(\frac{1}{z}\right)=O\left(\frac{1}{z}\right)$ and

$$
A^{(\infty)}= \begin{cases}z^{r_{\infty}-1}\left(A_{0}^{(\infty)}+\frac{A_{1}^{(\infty)}}{z}+\cdots+\frac{A_{r^{\infty}-1}^{(\infty)}}{z^{r \infty}-1}\right) & r_{\infty} \geq 1 \\ 0 & r_{\infty} \leq 0\end{cases}
$$

Lemma 4.4. $A(z)$ is rational.
Proof. $A(z)-\sum A^{(j)}(z)-A^{(\infty)}(z)$ is analytic on $\overline{\mathbb{C}}$ and tends to 0 as $z \rightarrow \infty$, so by Liouville is constant.

Definition (Poincaré rank). $r_{j}$ 's and $r_{\infty}$ are called Poincaré ranks of the singularities.

### 4.6 Holomorphic and meromorphic equivalence

Suppose $G: U(a) \backslash\{a\} \rightarrow \mathrm{GL}(n, \mathbb{C})$ analytic with $z=a$ at most a pole. Consider a system $\frac{d Y}{d z}=A(z) Y$. Given a gauge transformation $Y=G(z) \tilde{Y}$,

$$
\frac{d \tilde{Y}}{d z}=\underbrace{\left(G(z)^{-1} A(z) G(z)-G(z)^{-1} \frac{d G(z)}{d z}\right)}_{\tilde{A}(z)} \tilde{Y}
$$

(we assume $G(z)$ is invertible away from $z=a$ ).
Definition (holomorphic/meromorphic equivalence). We say two systems $\frac{d Y}{d z}=A(z) Y, \frac{d \tilde{Y}}{d z}=\tilde{A}(z) Y$, where $A, \tilde{A}$ are analytic on $U(a) \backslash\{a\}$ with at most a pole at $z=a$, are holomorphic equivalent (resp. meromorphic equivalent) at $z=a$ if exists $G(z)$ invertible in $U(a)$ (resp. $U(a) \backslash\{a\})$ which transforms one system into the other.

One checks that this is an equivalence relation.
Theorem 4.5. If $z=a$ is a regular singularity of $\frac{d Y}{d z}=A(z) Y$ then it is a meromorphic equivalent to a Fuchsian system $\frac{d \tilde{Y}}{d z}=\frac{L}{z-a} \tilde{Y}$.

Proof. Relabel the solution

$$
\mathbb{Y}(\tilde{z})=\underbrace{\Phi(z)}_{G(z)} \underbrace{(\tilde{z}-a)^{L}}_{\tilde{\mathbb{Y}}(z)}
$$

and $\tilde{\mathbb{Y}}(z)$ is a fundamental solution of the required system.
We can also consider formal equivalence. Consider the formal series

$$
\begin{aligned}
& A(z)=\frac{1}{(z-a)^{p+1}} \sum_{k=0}^{\infty} A_{k}(z-a)^{k} \\
& \tilde{A}(z)=\frac{1}{(z-a)^{q+1}} \sum_{k=0}^{\infty} \tilde{A}_{k}(z-a)^{k} \\
& G(z)=\frac{1}{(z-a)^{m}} \sum_{k=0}^{\infty} G_{k}(z-a)^{k}
\end{aligned}
$$

where $\operatorname{det} G_{0} \neq 0$. If formally $A$ and $\tilde{A}$ are related by the equation for gauge transformation above then the systems are formally meromorphic equivalent.

Remark. $G$ has a formal inverse

$$
\begin{aligned}
G(z)^{-1} & =(z-a)^{m}\left(G_{0}+\sum_{k=1}^{\infty} G_{k}(z-k)^{k}\right)^{-1} \\
& =(z-a)^{m}\left(I+\sum_{k=1}^{\infty} G_{0}^{-1} G_{k}(z-a)^{k}\right)^{-1} G_{0}^{-1} \\
& =(z-a)^{m}\left(I+\sum_{n=1}^{\infty}(-1)^{n}\left(\sum_{k=1}^{\infty} G_{0}^{-1} G_{k}(z-a)^{k}\right)^{n}\right) G_{0}^{-1}
\end{aligned}
$$

## 5 Structure offundamental matrix at a Fuchsian singularity

In this chapter we study the local structure of fundamental solution by doing example computations. Throughout consider the system $\frac{d Y}{d z}=\frac{A(z)}{z} Y$ where $A$ holomorphic at $z=0$ and $A(0) \neq 0$. Write $A(z)=A_{0}+\sum_{i=1}^{\infty} A_{i} z^{i}$. Let $J=G_{0}^{-1} A_{0} G_{0}$ be the Jordan normal form.

The constant gauge transformation $Y(z)=G_{0} \hat{Y}(z)$ gives

$$
\frac{d \hat{Y}}{d z}=\frac{G_{0}^{-1} A(z) G_{0}}{z} \hat{Y}
$$

so

$$
\hat{A}(z)=J+\sum_{i=1}^{\infty}\left(G_{0}^{-1} A_{i} G_{0}\right) z^{i}
$$

Suppose we have put $A_{0}$ into Jordan normal form, i.e. $\frac{d Y}{d z}=\frac{A(z)}{z} Y, A(z)=$ $J+\sum_{i=1}^{\infty} A_{i} z^{i}=\sum_{i=0}^{\infty} A_{i} z^{i}$. We seek a gauge transformation $G(z)$ such that $\frac{d \tilde{Y}}{d z}=\frac{R(z)}{z} \tilde{Y}$ where $R(z)=\sum_{i=0}^{\infty} R_{i} z^{i}$ is "as simple as possible" so that we can solve for $\tilde{Y}$.

The condition can be expressed as the expression

$$
G^{\prime}(z) \tilde{Y}+\underbrace{G(z) \tilde{Y}^{\prime}}_{\frac{R(z)}{z} \tilde{Y}}=\frac{A(z)}{z} G(z) \tilde{Y}
$$

so we aim to "solve"

$$
\begin{equation*}
z G^{\prime}+G R=A(z) G \tag{*}
\end{equation*}
$$

a system for $(G(z), R(z))$. Formally this means we are looking for series $G(z)=$ $\sum_{i=0}^{\infty} G_{i} z^{i}, A(z)=\sum_{i=0}^{\infty} A_{i} z^{i}, R(z)=\sum_{i=0}^{\infty} R_{i} z^{i}$ that satisfy ( $*$ ). After some manipulations we get

$$
\sum_{\ell=0}^{\infty}\left(\ell G_{\ell}+\sum_{j=0}^{\ell} G_{j} R_{\ell-j}\right) z^{\ell}=\sum_{\ell=0}^{\infty}\left(\sum_{j=0}^{\ell} A_{\ell-j} G_{j}\right) z^{\ell}
$$

so for $\ell=0$ we have $G_{0} R_{0}=A_{0} G_{0}$., i.e. $G_{0} R_{0}=J G_{0}$. We choose $G_{0}=I, R_{0}=$ $J$. For $\ell \geq 1$ we have

$$
G_{\ell}(\ell \cdot I+J)-J G_{\ell}=\left(\sum_{j=1}^{\ell-1}\left(A_{\ell-j} G_{j}-G_{j} R_{\ell-j}\right)+A_{\ell}\right)-R_{\ell}
$$

for $\left(G_{\ell}, R_{\ell}\right)$.

### 5.1 Case of diagonalisable $A_{0}$

Theorem 5.1. If $J=\Lambda$ then the system $\frac{d Y}{d z}=\frac{A(z)}{z} Y$ is holomorphically equivalent through $G_{0} G(z)$ to

$$
\frac{d \tilde{Y}}{d z}=\frac{1}{z}\left(\Lambda+R_{1} z+\cdots+R_{k} z^{k}\right) \tilde{Y}
$$

where $\left(R_{\ell}\right)_{i j} \neq 0$ only if $\lambda_{i}-\lambda_{j}=\ell \geq 1$ integer. This is called the normal form of the system. In particular if no difference $\lambda_{i}-\lambda_{j}$ is a non-zero integer then $\frac{d \tilde{Y}}{d z}=\frac{\Lambda}{z} \tilde{Y}$.

Proof. We show formal equivalence only and defer convergence to Theorem 5.5. Equation ( $\dagger$ ) reads for $\ell=1$

$$
G_{1} \Lambda-\Lambda G_{1}+G_{1}=A_{1}-R_{1}
$$

Taking $(i, j)$ component,

$$
\left(G_{1}\right)_{i j}\left(\lambda_{j}-\lambda_{i}+1\right)=\left(A_{1}\right)_{i j}-\left(R_{1}\right)_{i j}
$$

so if $\lambda_{i}-\lambda_{j} \neq 1$ then choose $R_{1}=0$ and $\left(G_{1}\right)_{i j}=\frac{\left(A_{1}\right)_{i j}}{\lambda_{j}-\lambda_{i}+1}$. If $\lambda_{i}-\lambda_{j}=1$ then $\left(R_{1}\right)_{i j}=\left(A_{1}\right)_{i j}$ and $\left(G_{1}\right)_{i j}$ is arbitrary.

Similarly for $\ell \geq 1$ in general

$$
G_{\ell} \Lambda-\Lambda G_{\ell}+\ell G_{\ell}=(\cdots)-R_{\ell}
$$

so

$$
\left(G_{\ell}\right)_{i j}\left(\lambda_{j}-\lambda_{i}+\ell\right)=(\cdots)_{i j}-\left(R_{\ell}\right)_{i j}
$$

If $\lambda_{i}-\lambda_{j} \neq \ell$ then choose $\left(R_{\ell}\right)_{i j}=0$ and $\left(G_{\ell}\right)_{i j}=\frac{(\cdots)}{\lambda_{j}-\lambda_{i}+\ell}$. If $\lambda_{i}-\lambda_{j}=\ell$ then we choose $\left(R_{\ell}\right)_{i j}=(\cdots)_{i j}$ and $\left(G_{\ell}\right)_{i j}$ is arbitrary.

Since there are only finitely many $\lambda_{i}-\lambda_{j}$, the system is of the form required.

Corollary 5.2. The system, if $A_{0}$ is diagonalisable, has a fundamental solution

$$
Y(z)=G_{0} G(z) z^{\Lambda} z^{R}
$$

where $R=R_{1}+\cdots+R_{k}$ (which is nilpotent). In particular $z=0$ is a Fuchsian singularity so regular (because $\operatorname{det} G(z) \neq 0$ ).
Proof. We know $Y(z)=G_{0} G(z) \tilde{Y}(z)$ where $\tilde{Y}(z)$ satisfies ( $\ddagger$ ). Claim that $\tilde{Y}(z)=z^{\Lambda} z^{R}$ is a fundamental solution:

$$
\frac{d}{d z}\left(z^{\Lambda} z^{R}\right)=\frac{\Lambda}{z} z^{\Lambda} z^{R}+z^{\Lambda} \frac{R}{z} z^{R}=\frac{1}{z}\left(\Lambda+z^{\Lambda} R z^{-\Lambda}\right) z^{\Lambda} z^{R}
$$

and

$$
\left(z^{\Lambda} R z^{-\Lambda}\right)_{i j}=z^{\lambda_{i}-\lambda_{j}} R_{i j}= \begin{cases}z^{\ell}\left(R_{\ell}\right)_{i j} & \lambda_{i}-\lambda_{j}=\ell \geq 1 \\ 0 & \lambda_{i}-\lambda_{j} \neq \ell \geq 1\end{cases}
$$

so indeed it is a fundamental solution.

Definition (resonance). If $\lambda_{i}-\lambda_{j} \in \mathbb{Z} \backslash\{0\}$ we say the system is resonant at $z=0$.

Exercise. Use the same strategy to prove

$$
e^{2 \pi i \Lambda} R=R e^{2 \pi i \Lambda}
$$

which implies that

$$
e^{2 \pi i \Lambda} z^{R}=z^{R} e^{2 \pi i \Lambda}
$$

The monodromy of the solution is given by

$$
Y(z)=G_{0} G(z) z^{\Lambda} z^{R} \mapsto G_{0} G(z) z^{\Lambda} e^{2 \pi i \Lambda} z^{R} e^{2 \pi i R}=\underbrace{G_{0} G(z) z^{\Lambda} z^{R}}_{Y(z)} \underbrace{e^{2 \pi i \Lambda} e^{2 \pi i R}}_{M} .
$$

We can reduce this further: write $\lambda_{j}=d_{j}+\rho_{j}$ where $d_{j} \in \mathbb{Z}, 0 \leq \operatorname{Re} \rho_{j}<1$ so $\Lambda=D+S$ where both $D$ and $S$ and diagonal and $D$ has integer entries. Then $[D, S]=0$. Also $[R, S]=0$ since

$$
[R, S]_{i j}=\left(\rho_{j}-\rho_{i}\right) R_{i j}
$$

so either $\lambda_{i}-\lambda_{i} \in \mathbb{Z}$ so $\rho_{j}-\rho_{i}=0$, or $\lambda_{i}-\lambda_{j} \notin \mathbb{Z}$ so $R_{i j}=0$. Thus

$$
Y(z)=G_{0} G(z) z^{D+S} z^{R}=G_{0} G(z) z^{D} z^{S} z^{R}=G_{0} G(z) z^{D} z^{R+S}=\underbrace{G_{0} G(z) z^{D}}_{\Phi(z)} z^{L}
$$

where $L=R+S$. But this has exactly the form $Y(z)=\Phi(z) \cdot z^{L}$ where $\Phi=G_{0} G(z) z^{D}$ is single valued with a pole at $z=0$. In this representation the monodromy matrix is simply $M=e^{2 \pi i L}$.

Let's discuss for a moment the freedom in the solution.

1. In $G_{0}^{-1} A_{0} G_{0}=J, G_{0}$ has freedom $G_{0} \mapsto G_{0} \Delta_{0}$ where $\Delta_{0}^{-1} J \Delta_{0}=J$.
2. Fix $G_{0}$ and $\Lambda$. Recall that $G_{\ell}$ contains a finite number of arbitrary parameters.
3. Also $R$ is fixed. Suppose

$$
Y(z)=G_{0}\left(I+\sum G_{j} z^{j}\right) z^{D} z^{L} .
$$

Suppose we can find
4. Freedom in $R$. The system may be put in two different normal forms with same $\Lambda$. Two such solutions $\tilde{Y}$ and $\tilde{Y}$ are related by a gauge transformation $\tilde{Y}=\Delta(z) \tilde{Y}$.

$$
\Delta(z)=\Delta_{0}\left(I+\Delta_{1} z+\cdots+\Delta_{k} z^{k}\right)
$$

where $\left(\Delta_{\ell}\right)_{i j} \neq 0$ only if $\lambda_{i}-\lambda_{j} \ell \geq 1$ and i

### 5.2 Case of general $A_{0}$

In general, any $A_{0}$ can be written in Jordan normal form $J=\operatorname{diag}\left(J_{1}, \ldots, J_{s}\right)$ where $J_{\ell}-\lambda_{\ell} I_{\ell}+H_{\ell}$ where $H_{\ell}$ has only entries 1 right above the diagonal. We can then partition $R_{\ell}$ into blocks and the the problem effectively becomes solving the system of equations

$$
X A-B X=C
$$

where $A$ is $m \times m, X$ is $n \times m, B$ is $n \times n$ and $C$ is $m \times m$. Then

1. $X A-B X=0$ has nontrivial solution if and only if $A, B$ have a common eigenvalue.
2. $X A-B X=C$ with $C$ given has a unique solution $X$ if and only if $A, B$ hasve non common eigenvalues.
3. $X A-B X=C$, with both $X, C$ unknown and $A, B$ with at least one common eigenvalue, always has a solution.

All these can be proven by regarding $X$ and $C$ as vectors and consider the linear system $D x=c$. Then $\operatorname{det} D \neq 0$ if and only if $A, B$ have no common eigenvalues. See Wasow.

Subsequently we can divide into cases according to if $\lambda_{i}-\lambda_{j}=\ell$. Thus we have the following more general theorem

Theorem 5.3. The system $\frac{d Y}{d z}=\frac{A(z)}{z} Y$ is holomorphically equivalent through $G_{0} G(z)$ to

$$
\frac{d \tilde{Y}}{d z}=\frac{1}{z}\left(J+R_{1} z+\cdots+R_{k} z^{k}\right) \tilde{Y}
$$

where $\left(R_{\ell}\right)_{i j} \neq 0$ only if $\lambda_{i}-\lambda_{j}=\ell \geq 1$ integer.
Let $\lambda_{j}=d_{j}+r_{j}$ as before and write $S=J-D=\operatorname{diag}\left(S_{1}, \ldots, S_{s}\right)$ where $S_{j}=\rho_{j} I_{j}+H_{j}$. Let $L=R+S$ where $R=R_{1}+\cdots+R_{r}$.

Corollary 5.4. The system $\frac{d Y}{d z}=\frac{A(z)}{z} Y$ has a fundamental solution

$$
Y(z)=G_{0} G(z) z^{D} z^{L}
$$

where $G(z)=I+\sum_{j=0}^{\infty} G_{j} z^{j}$.

### 5.3 Convergece of formal solution

We now prove covergence. Recall (*)

$$
z G^{\prime}+G R(z)=A(z) G
$$

assume $R(z)$ is a polynomial, $A(z)$ analytic at $z=0$. Construct $y=\left(G_{11}, \ldots, G_{n n}\right)^{T}$, (*) becomes

$$
z y^{\prime}=F(z) y
$$

where $F(z)$ is analytic in a ball $U_{r}(0)$ around $z=0$ so can be locally written as a convergent power series.

Theorem 5.5. Suppose $z y^{\prime}=F(z) y$ has a formal solution $y(z)=\sum c_{j} z^{j}$. Then the series converges locally uniformly on $U_{r}(0)$ and thus defines an analytic function in $U_{r}(0)$.

Note this is really a property of Fuchsian singularity as if we change LHS to $z^{p} y^{\prime}$ where $p>1$ then the conclusion does not hold.

Proof. Recall from complex analysis that if $f(z)$ is holomorphic on $U_{r}(0)$ then $f(z)=\sum c_{k} z^{k}$ where

$$
c_{k}=\frac{f^{(k)}(0)}{k!}=\frac{1}{2 \pi i} \oint_{C_{R_{1}}} \frac{f(\zeta)}{\zeta^{k+1}} d \zeta
$$

so we can estimate

$$
\left|c_{k}\right| \leq \frac{1}{2 \pi} \oint_{C_{R_{1}}} \frac{\mid f(\zeta)}{|f|^{k+1}} d|\zeta| \leq \frac{M\left(R_{1}\right)}{R_{1}^{k}}
$$

Conversely given $\sum c_{k} z^{k}$, assuming exists $M$ such that $\left|c_{k}\right| \leq \frac{M}{R_{1}^{k}}$, then the series is uniformly convergence on $|z| \leq R_{2}$ for all $R_{2} \leq R_{1}$, i.e. it defines an analytic function on $U_{R_{1}}(0)$.

Now given a formal solution $y=\sum c_{j} z^{j}$, substitute into the equation to get

$$
\sum j C_{j} z^{j}=\sum_{i j} F_{j} C_{j} z^{i+j}
$$

so

$$
\sum \ell C_{\ell} z^{\ell}=\sum_{\ell}\left(\sum_{j=0}^{\ell} F_{\ell-j} C_{j}\right) z^{\ell}
$$

SO

$$
\left(\ell I-F_{0}\right) C_{\ell}=F_{\ell} C_{0}+F_{\ell-1} C_{1}+\cdots+F_{1} C_{\ell-1}
$$

For $\ell \geq N$ sufficiently large such that $\frac{\left|F_{0}\right|}{\ell} \leq 1$, then $\left(\ell I-F_{0}\right)$ has an inverse and we can solve for $C_{\ell}$. We can esimate

$$
\begin{aligned}
\left|C_{\ell}\right| & \leq \frac{1}{\ell} \sum_{k=0}^{\infty} \frac{\left|F_{0}\right|^{k}}{\ell^{k}}\left(\left|F_{\ell}\right|\left|C_{0}\right|+\cdots+\left|F_{1}\right| \mid C_{\ell-1 \mid}\right) \\
& =\frac{1}{\ell} \frac{1}{1-\left|f_{0}\right| / \ell}\left(\left|F_{\ell}\right|\left|C_{0}\right|+\cdots+\left|F_{1}\right| \mid C_{\ell-1 \mid}\right) \\
& \leq \frac{C}{\ell}\left(\left|F_{\ell}\right|\left|C_{0}\right|+\cdots+\left|F_{1}\right| \mid C_{\ell-1 \mid}\right)
\end{aligned}
$$

We know that $F$ is analytic in $U_{r}(0)$ so $\left|F_{k}\right| \leq \frac{M}{R_{1}^{k}}$. Thus

$$
\left|C_{\ell}\right| \leq \frac{M C}{\ell}\left(\frac{\left|C_{0}\right|}{R_{1}^{\ell}}+\cdots+\frac{C_{\ell-1}}{R_{1}}\right) \leq M C\left(\frac{\left|C_{0}\right|}{R_{1}^{\ell}}+\cdots+\frac{C_{\ell-1}}{R_{1}}\right)
$$

If $\ell \leq N-1(?)$ we can always choose $P>0$ such that $\left|C_{\ell}\right| \leq\left(\frac{R}{R_{1}}\right)^{\ell}$. Prove that for all $\ell, \mid C_{\ell} \leq\left(P / R_{1}\right)^{\ell}$ : assume this is true up to $\ell-1$ and prove it for $\ell$.

$$
\left|\mathbb{C}_{\ell}\right| \leq M C\left(\frac{\left|C_{0}\right|}{R_{1}^{\ell}}+\cdots+\frac{C_{\ell-1}}{R_{1}}\right) \leq \frac{M C}{R_{1}^{\ell}}\left(P^{\ell-1}+P^{\ell-2}+\cdots+1\right)=\frac{M C}{R_{1}^{\ell}} \frac{P^{\ell}-1}{P-1}
$$

$P^{\ell-1}(P-1-M C) \geq 0$ this is possible, provided that $P$ is sufficiently large.
Thus $\left|C_{\ell}\right| \leq \ldots$ so $\sum\left|C_{\ell}\right||z|^{\ell}<\infty$ provided $|z| \leq \frac{R_{1}}{P}$, so $y(z)$ is an analytic solution for $|z| \leq \frac{R_{1}}{P}$. From general theory we can analytically continue $y(z)$ as analytic solution in ball $|z|<r$.

Thus every time we have a Fuchsian singularity, all formal computation are actual (meaning analytic) computation.

### 5.4 Completely Fuchsian system

We have studied the local theory of a Fuchsian system. In this section we study the behaviour if the a general system in which all the singularities in the complex plane are Fuchsian. Recall $A(z)$ is rational so necessarily such a system has the form

$$
\frac{d Y}{d z}=\left(\frac{A_{1}}{z-a_{1}}+\cdots+\frac{A_{m}}{z-a_{m}}\right) Y
$$

where $A_{j}$ 's are constant matrices. At $z=\infty$,

$$
A(z)=\frac{1}{z} \sum_{j=1}^{m} A_{j}+O\left(\frac{1}{z^{2}}\right)
$$

which is convergent for $|z|>R$. Substitute $z=\frac{1}{t}$, the system becomes

$$
\frac{d Y}{d t}=-\frac{1}{t^{2}} A\left(\frac{1}{t}\right) Y=\frac{1}{t}(\underbrace{-\sum A_{j}}_{A_{\infty}}+O(t)) Y
$$

For each singularity $a_{j}$, we know from local theory that there exists a neighbourhood $U\left(a_{j}\right)$ on which

$$
Y_{j}(z)=G_{j}(z)\left(z-a_{j}\right)^{D_{j}}\left(z-a_{j}\right)^{L_{j}} .
$$

Suppose $A_{j}$ has eigenvalues $\lambda_{1}^{(j)}, \cdots \lambda_{m}^{(j)}$ and $\lambda_{k}^{(j)}=d_{k}^{(j)}+\rho_{k}^{(j)}$. Similar for $z=\infty$ : in a neighbourhood of $t=0$,

$$
Y_{\infty}(z)=Y(t)=\tilde{G}_{\infty}(t) t^{D_{\infty}} t^{L_{\infty}}=G_{\infty}(z) z^{-D_{\infty}} z^{-L_{\infty}}
$$

where $G_{\infty}(z)$ is analytic at $z=\infty$.
Remark. To talk about the local solutions we must select a sheet of the universal cover. It is given as follow. Choose parallel branch cuts $L_{j}$ which are rays from $a_{j}$ 's in direction $\eta$, such that no $a_{k}$ is contained $L_{j}$ for $j \neq k$. Then we stipulate that either $\eta-2 \pi<\arg \left(z-a_{j}\right)<\eta$ for $\left|z-a_{j}\right|$ small, or $\eta-2 \pi<\arg z<\eta$ for $|z|$ large.

Remark. Note that from definition $A_{\infty}+\sum_{j=1}^{m} A_{j}=0$. Taking the trace, we get

$$
\sum_{k=1}^{n}\left(\lambda_{k}^{(1)}+\lambda_{k}^{(2)}+\cdots+\lambda_{k}^{(m)}+\lambda_{k}^{(\infty)}\right)=0
$$

This interdependence of eigenvalues is called Fuchs identity.
Since each one of $Y_{j}(z)$ and $Y_{\infty}(z)$ is a fundamental solution, we can relate them by $Y_{\infty}(z)=Y_{j}(z) C_{j}$ where $C_{j}$ nonsingular is called the connection matrix. For computation, $G_{j}(z), G_{\infty}(z), D_{j}, L_{j}, D_{\infty}, L_{\infty}$ are algebraic functions of $A_{1}, \ldots, A_{m}$ so we can compute them. However $C_{j}$ 's are transcendental functions of $A_{i}$ 's.

Note $C_{j}=Y_{j}\left(z_{0}\right)^{-1} Y_{\infty}\left(z_{0}\right)$ for $z_{0} \notin\left\{a_{1}, \ldots, a_{m}, \infty\right\}$ and we know each $Y_{j}(z)$ only locally so the study of the global Fuchsian system is reduced to asking how to compute $Y_{j}\left(z_{0}\right), Y_{\infty}\left(z_{0}\right)$ when $z_{0} \notin U\left(a_{1}\right)$ or $z_{0} \notin U(\infty)$. This is called the connection problem.

In several cases $(n=2)$ this can be solved (using linear special function theory such as Bessel, Airy, hypergeometric functions). The general strategy is as follow: for the $(a, b)$ entry we write

$$
\left(Y_{j}(z)\right)_{a b}=\int_{\gamma} \phi_{a b}(z, s) d s
$$

called integral representation. Then

$$
\left(Y_{j}\left(z_{0}\right)\right)_{a b}=\int_{\gamma} \phi_{a b}\left(z_{0}, s\right) d s
$$

which is computable in terms of classical special functions.

Monodromy of global Fuchsian system We know the monodromy of the local solutions: for $\left|z-a_{j}\right|$ small, if we take the loop $\gamma_{j}$ that goes around $a_{j}$ counterclockwise once we have

$$
Y_{j}(z) \mapsto Y_{j}(z) e^{2 \pi i L_{j}} .
$$

For $Y_{\infty}$, we take big $|z|$ and loop around $\infty$ clockwise to get

$$
Y_{\infty}(z) \mapsto Y_{\infty}(z) e^{2 \pi i L_{\infty}} .
$$

Combining the local monodromy with connection matrices,

$$
Y_{\infty}(z)=Y_{j}(z) C_{j} \mapsto_{\gamma_{j}} Y_{j}(z) e^{2 \pi i L_{j}} C_{j}=Y_{\infty}(z) \underbrace{C_{j}^{-1} e^{2 \pi i L_{j}} C_{j}}_{M_{j}}
$$

so the monodromy (anti)representation is given by

$$
\gamma_{j} \mapsto M_{j}=C_{j}^{-1} e^{2 \pi i L_{j}} C_{j}
$$

Note that $\gamma_{1} \cdots \gamma_{2} \cdots \gamma_{m}=\gamma_{\infty}^{-1}$ so $M_{\infty}=\left(M_{m} \cdots M_{1}\right)^{-1}$. Again we stress that the difficulty lies in finding $C_{j}$.

## 6 Linear equations of order $n$

Consider the scalar ODE of order $n$

$$
u^{(n)}+a_{1}(z) u^{(n-1)}+\cdots+a_{n}(z) u=0
$$

where $z=0$ is an isolated singularity of $a_{j}(z) . a_{j}(z)$ are analytic the we can express the solutions in Taylor series. For the simplest case $a_{j} \in \mathbb{C}$, we can solve the indicial equation

$$
f_{0}(\lambda)=\lambda^{n}+a_{1} \lambda^{n-1}+\cdots+a_{n}=\prod_{j=1}^{s}\left(\lambda-\lambda_{j}\right)^{m_{j}}
$$

and $e^{\lambda_{j} z}, z e^{\lambda_{j} z}, \ldots, z^{m_{j}-1} e^{\lambda_{j} z}$ for $j=1, \ldots, s$ is a fundamental system.
As another example

$$
u^{(n)}+\frac{a_{1}}{z} u^{(n-1)}+\cdots+\frac{a_{n}}{z^{n}} u=0
$$

where $a_{j} \in \mathbb{C}$. This is called the Euler equation. For $z \neq 0$ multiply by $z^{n}$ to get

$$
E(u)=0
$$

where $E=z^{n} \frac{d^{n}}{d z^{n}}+\cdots+a_{n}$. Define the Euler operator $\delta=z \frac{d}{d z}$ and define

$$
[\delta]_{k}=\delta(\delta-1) \cdots(\delta-k+1)
$$

and a quick calculation shows $[\delta]_{k}=z^{k} \frac{d^{k}}{d z^{k}}$ so

$$
E=[\delta]_{n}+a_{1}[\delta]_{n-1}+\cdots+a_{2} \delta+a_{1}
$$

and we have another indicial equation

$$
f_{0}(\lambda)=[\lambda]_{n}+a_{1}[\lambda]_{n-1}+\cdots+a_{n}=\prod_{j=1}^{s}\left(\lambda-\lambda_{j}\right)^{m_{j}}
$$

where $\lambda_{i}$ 's are distinct. They are called indices or characteristic exponents.

Frobenius method A computation shows

$$
[\delta]_{k} z^{\lambda}=[\lambda]_{k} z^{\lambda}
$$

so $z^{\lambda}$ is an eigenfunction so

$$
E\left(z^{\lambda}\right)=f_{0}(\lambda) z^{\lambda} .
$$

Thus if $\lambda_{i}$ 's are the solutions of the indicial equation then $z^{\lambda_{i}}$ 's are solutions to the Euler equation.

Since

$$
\frac{\partial^{\ell}}{\partial \lambda^{\ell}} z^{\lambda}=(\log z)^{\ell} z^{\lambda}
$$

it follows that

$$
E\left(\frac{\partial^{\ell} z^{\lambda}}{\partial \lambda^{\ell}}\right)=\frac{\partial^{\ell}}{\partial \lambda^{\ell}} E\left(z^{\ell}\right)=\frac{\partial^{\ell}}{\partial \lambda^{\ell}}\left(f_{0}(\lambda) z^{\lambda}\right)=\sum_{k=0}^{\ell}\binom{\ell}{k} \frac{d^{k} f_{0}(\lambda)}{d \lambda^{k}}(\log z)^{\ell-k} z^{\lambda} .
$$

If $\lambda_{j}$ is a root of $f_{0}$ with multiplicity $m_{j}$ then $E\left(\frac{\partial^{\ell} z^{\ell}}{\partial \lambda^{\ell}}\right)=0$ for $\ell<m_{j}$ so

$$
z^{\lambda_{j}}, \ln z \cdot z^{\lambda_{1}}, \ldots, z^{\lambda_{j}}(\ln z)^{m_{j}-1}
$$

form a fundamental system.
In this example the companion system is

$$
\frac{d Y}{d z}=\left(\begin{array}{ccc}
0 & 1 & \\
& 0 & 1 \\
\vdots & & \\
-\frac{a_{m}}{z^{m}} & & -\frac{a_{1}}{z}
\end{array}\right) y
$$

for which $z=0$ is a singularity of second kind if $m \geq 2$. In addition we know this is a regular singularity from the fundamental system above. It is not clear, however, that this system is a priori regular.

We can make this more lucid by making the substituion

$$
y_{1}=u, y_{2}=z u^{\prime}, \ldots, y_{n}=z^{n-1} u^{(n-1)}
$$

so

$$
y_{j}^{\prime}= \begin{cases}\frac{1}{z}\left((j-1) y_{j}+y_{j+1}\right) & j \leq n-1 \\ \frac{1}{z}\left((n-1) y_{n}-\left(a_{1} y_{n}+\cdots+a_{n} y_{1}\right)\right) & j=n\end{cases}
$$

so the system is

$$
\frac{d Y}{d z}=\frac{1}{z}\left(\begin{array}{ccccc}
0 & 1 & & & \\
& 1 & 1 & & \\
\vdots & & 2 & 1 & \\
-a_{n} & -a_{n-1} & \cdots & -a_{2} & n-1-a
\end{array}\right) y
$$

and it is obvious that $z=0$ is a Fuchsian singularity.
Call the constant matrix $A_{0}$ so from general theory of Fuchsian singularity

$$
Y(z)=z^{A_{0}}=G_{0} z^{J} G_{0}^{-1} .
$$

One can also check $\operatorname{det}\left(\lambda-A_{0}\right)=f_{0}(\lambda)$. Note that the original system has Poincaré rank $n-1$, while after sustituion it has Poincaré rank 0 . The solutions are related by a meromorphic tranformation

$$
\left(\begin{array}{ccccc}
1 & & & & \\
& \frac{1}{z} & & & \\
& & \frac{1}{z^{2}} & & \\
& & & \ddots & \\
& & & & \frac{1}{z^{n-1}}
\end{array}\right)
$$

This is called a shearing transformation.

### 6.1 Series solution

Consider

$$
u^{(n)}+\frac{\hat{a}_{1}(z)}{z} u^{(n-1)}+\cdots+\frac{\hat{a}_{n}(z)}{z^{n}}(u)=0
$$

where $\hat{a}_{j}(z)$ analytic at $z=0$. The only difference between this and the Euler equation is that the coefficients are not constant. Write $E(u)=0$. Use the Frobenius method with a pertubation by substituting

$$
u(z)=z^{\lambda} \sum_{\nu=0}^{\infty} c_{\nu} z^{\nu}
$$

$E\left(\sum_{\nu=0}^{\infty} c_{\nu} z^{\nu+\lambda}\right)=\sum_{\nu=0}^{\infty} c_{\nu} E\left(z^{\nu+\lambda}\right)=\sum_{\nu=0}^{\infty} c_{\nu} \underbrace{\left([\lambda+\nu]_{n}+\hat{a}_{1}(z)[\nu+\lambda]_{n-1}+\cdots+\hat{a}_{n}(z)\right)}_{\bullet} z^{\nu}$
Substitute $\hat{a}_{j}(z)=\sum_{\ell=0}^{\infty} a_{\ell}^{(j)} z^{\ell}$, we can write ( $\bullet$ ) as

$$
\sum_{\ell=0}^{\infty} f_{\ell}(\lambda+\nu) z^{\ell}
$$

for some $f_{\ell}$. Thus
$E\left(\sum_{\nu} c_{\nu} z^{\nu+\lambda}\right)=\sum_{\nu, \ell} c_{\nu} f_{\ell}(\lambda+\nu) z^{\nu+\ell}=\sum_{m=0}\left(\sum_{\nu=0}^{m} c_{\nu} f_{m-\nu}(\lambda+\nu)\right) z^{m} \quad m=\nu+\ell$
so for $E=0$ we reuire all coefficients of $z^{m}$ to vanish, so we get recurrence relations

$$
\begin{aligned}
c_{0} f_{0}(\lambda) & =0 \\
c_{1} f_{0}(\lambda+1)+c_{0} f_{1}(\lambda) & =0 \\
c_{m} f_{0}(\lambda+m)+\cdots+c_{0} f_{m}(\lambda) & =c
\end{aligned}
$$

Write $f_{0}(\lambda)=\prod\left(\lambda-\lambda_{j}\right)^{m_{j}}$. If $\lambda=\lambda_{j}$ then $c_{0}$ is arbitrary. If in addition $\lambda_{j}+m$ is not a root then we can determine all $c_{m}$ in terms of $c_{0}$ and $\lambda_{j}$.

$$
u(z)=z^{\lambda_{j}}\left(c_{0}+\sum_{\nu=1}^{\infty} c_{\nu}\left(c_{0}, \lambda_{j}\right) z^{\nu}\right)
$$

In summary, if $f_{0}(\lambda)$ has $n$ distinct roots $\lambda_{1}, \ldots, \lambda_{n}$ not differing by integers then we have a fundamental system

$$
u(z)_{j}=z^{\lambda_{j}}\left(c_{0}^{(j)}+\sum_{\nu=1}^{\infty} c_{\nu}^{(j)} z^{\nu}\right)
$$

If some $m_{j} \geq 2$, i.e. multiple roots, or if $\lambda_{j}+m$ is also a root, we can implement Frobenius method to get $z^{\lambda_{j}}(\ln z)^{k}$. For reference see Ince.

A few more words on where $\ln z$ comes from

Similar as before we can reduce the ODE to a linear system by subsituting $y_{1}=u, y_{2}=z u^{\prime}, \ldots, y_{n}=z^{n-1} u^{(n-1)}$, and consider the system

$$
\frac{d y}{d z}=\frac{1}{z}\left(\begin{array}{cccccc}
0 & 1 & & & \\
& 1 & 1 & & \\
& & 2 & 1 & \\
& & & \ddots & \\
-\operatorname{hata}_{n}(z) & -\hat{a}_{n-1}(z) & \cdots & -\hat{a}_{1}(z)+n-1 & & \\
& & & \\
& & & \\
& & &
\end{array}\right) y
$$

and it is obvious that $z=0$ is a Fuchsian singularity, ergo the name. Since $\hat{a}(z)$ is analytic at $z=0$, write $\hat{a}(z)=\sum_{\ell=0}^{\infty} a_{\ell}^{(j)} z^{\ell}$. Then we find

$$
A_{0}=A(0)=\left(\begin{array}{cccccc}
0 & 1 & & & & \\
& 1 & 1 & & \\
& & 2 & 1 & \\
& & & \ddots & & \\
-a_{0}^{(n)} & -a_{0}^{(n-1)} & \cdots & -a_{0}^{(1)}+n-1 & & 1
\end{array}\right)
$$

and we can $\operatorname{check} \operatorname{det}\left(\lambda-A_{0}\right)=f_{0}(\lambda)$.
We know that exists fundamental solution $Y(z)=G(z) z^{D} z^{S+R}$ where $G$ is holomorphically invertible at $z=0, D+S=J$, the Jordan normal form of $A_{0}$, and $R$ nilpotent. Suppose the first row of $Y(z)$ is $\left(u_{1}(z), \ldots, u_{n}(z)\right)$, which are solutions to the ODE. Then $u_{j}$ containes terms $z^{\lambda_{k}}(\ln z)^{p}$ where $1 \leq k<s$ and $p \leq m_{k}-1$.

Remark. The solutions obtained by Frobenius method converge because the companion system is Fuchsian. Thus the formal solutions are actual solutions.

A converse to proposition to Theorem 4.3:
Theorem 6.1. For an $O D E, z=0$ is Fuchsian if and only if it is a regular singularity.

Proof. Only need to prove only if. Induction on the order of the ODE $n$. For $n=1$, suppose $z=0$ is a regular singularity of $u^{\prime}+a_{1}(z) u=0$ and $u(z)$ is a solution. Then it has monodromy

$$
u\left(z e^{2 \pi i}\right)=\mu u(z)=e^{2 \pi i \lambda} u(z)
$$

where $\mu \in \mathbb{C} \backslash\{0\}$ and $\lambda$ is defined up to integers. Then $u(z) z^{-\lambda}$ is single valued so $z=0$ is regular, so $z=0$ is at most a pole. Choose $\lambda$ so that $f(z)=u(z) z^{-\lambda}$ is analytic at $z=0$ and $f(0) \neq 0$. Then

$$
a_{1}(z)=-\frac{u^{\prime}}{u}=-\frac{\lambda}{z}+\frac{f^{\prime}(z)}{f(z)}
$$

which has a pole of order 1 .
We do the induction step for $n=2$, consider $u^{\prime \prime}+a_{1}(z) u^{\prime}+a_{2}(z) u=0$ where $z=0$ is regular. Exists $u_{1}(z), u_{2}(z)$ linearly independent such that $\left(u_{1}, u_{2}\right)$ has
monodromy matrix $M$. By a linear transformation wlog $M$ is upper trianglular so $u_{1}\left(z e^{2 \pi i}\right)=\mu_{1} u_{1}(z)$ so as before $u_{1}(z)=f(z) z^{\lambda}$ where $f(z)$ analytic and $f(0) \neq 0$. We seek another solution $u(z)=u_{1}(z) v(z)$. Substitute the ansatz into the equation to get the condition

$$
v^{\prime \prime}+\left(\frac{2 u_{1}(z)^{\prime}}{u_{1}(z)}+a_{1}(z)\right) v^{\prime}=0
$$

a first order ODE in $v^{\prime}$. Since $z=0$ is regular, $v^{\prime}(z)$ has a regular singularity in $z=0$ so the first step applies to $v^{\prime}$ and conclude that

$$
\frac{2 u_{1}^{\prime}(z)}{u_{1}(z)}+a_{1}(z)=\frac{2 \lambda_{1}}{z}+\frac{2 f^{\prime}(z)}{f(z)}+a_{1}(z)
$$

has at most a simple pole at $z=0$. It remains to compute $a_{2}(z)$ :

$$
a_{2}(z)=-\left(\frac{u_{1}^{\prime \prime}}{u_{1}}+a_{1} \frac{u_{1}^{\prime}}{u_{1}}\right)=-\left(\frac{\lambda(\lambda-1)}{z^{2}}+\frac{\lambda}{z}\left(\frac{2 f^{\prime}}{f}+a_{1}(z)\right)+\frac{f^{\prime \prime}}{f}+a_{1}(z) \frac{f^{\prime}}{f}\right)
$$

so $a_{z}(z)$ has at most a pole of order 2 at $z=0$.
We can also define Fuchsian singularity at $\infty$ by settting $t=\frac{1}{z}$ and consider the behaviour at $t=0$.

Proposition 6.2. $z=\infty$ is Fuchsian if and only if $a_{j}(z)=\frac{\tilde{a}_{j}}{z^{j}}$ where $\tilde{a}_{j}(z)=\tilde{a}_{0}^{(j)}+\sum_{\ell=1}^{\infty} \frac{\tilde{a}_{\ell}^{(j)}}{z^{\ell}}$ is analytic at $z=\infty$.
Proof. $u^{(n)}+a_{1}(z) u^{(n-1)}+\cdots+a_{n}(z) u=0$. Put $a_{j}(z)=\frac{\tilde{a}_{j}(z)}{z^{j}}$ (we do not know the analytic property of $\left.\tilde{a}_{j}\right)$. Multiply by $z^{n}$ and substitute $z=\frac{1}{t}$, note

$$
\delta=z \frac{z}{d z}=-t \frac{d}{d t}=-\delta_{t}
$$

so $[\delta]_{k}=(-1)^{k}\left(\delta_{t}\right)_{k}$ where

$$
\left[\delta_{t}\right]_{k}=\delta_{t}\left(\delta_{t}+1\right) \cdots\left(\delta_{t}+n-1\right)
$$

Then the system is

$$
\left[\delta_{t}\right]_{n}-\tilde{a}_{1}\left(\frac{1}{t}\right)\left[\delta_{t}\right]_{n-1}+\cdots+(-1)^{n} \tilde{a}_{n}\left(\frac{1}{t}\right)
$$

so $t=0$ is Fuchsian if and only if $\tilde{a}_{j}\left(\frac{1}{t}\right)$ is analytic at $t=0$.

### 6.2 Completely Fuchsian ODE of order $n$

Suppose $z=\alpha_{1}, \ldots, \alpha_{m}, \infty$ are poles of the $a_{j}(z)$ 's, which have the correct order.

Proposition 6.3. The system is completely Fuchsian if and only if

$$
a_{j}(z)=\frac{p_{j}(z)}{\left(z-\alpha_{1}\right)^{j} \cdots\left(z-\alpha_{m}\right)^{j}}
$$

where $P_{j}$ is a polynomial of degree $\leq j \cdot m-j$.

Proof. At $z=\alpha_{k}$,

$$
a_{j}(z) \sim \frac{\operatorname{reg}\left(z-\alpha_{k}\right)}{\left(z-\alpha_{k}\right)^{j}}
$$

and at $z=\infty$,

$$
z^{j} a_{j}(z) \sim z^{j+\operatorname{deg} P_{j}-j \cdot m} .
$$

By partial fractions we can find the explicit forms of all completely Fuchsian system. As an example calculation, take $n=2$ and consider the equation

$$
u^{\prime \prime}+a_{1} u^{\prime}+a_{2} u=0
$$

Then

$$
\begin{aligned}
& a_{1}(z)=\frac{A_{1}}{z-\alpha_{1}}+\cdots+\frac{A_{m}}{z-\alpha_{m}} \\
& a_{2}(z)=\sum_{j=1}^{m}\left(\frac{B_{j}}{\left(z-\alpha_{j}\right)^{2}}+\frac{C_{j}}{z-\alpha_{j}}\right)
\end{aligned}
$$

where $A_{j}, B_{j}, C_{J} \in \mathbb{C}$ and we can show $\sum_{j=1}^{m} C_{j}=0$.

- For $m=1$ there is only one equation $u^{\prime \prime}=0$.
- For $m=2$,

$$
u^{\prime \prime}+\frac{A_{1}}{z-\alpha_{1}} u^{\prime}+\frac{B_{1}}{\left(z-\alpha_{1}\right)^{2}} u=0
$$

the Euler equation.

- For $m=3$,
$u^{\prime \prime}+\left(\frac{A_{1}}{z-\alpha_{1}}+\frac{A_{3}}{z-\alpha_{3}}\right) u^{\prime}+\left(\frac{B_{1}}{\left(z-\alpha_{1}\right)^{2}}+\frac{B_{2}}{\left(z-\alpha_{2}\right)^{2}}+C_{1}\left(\frac{1}{z-\alpha_{1}}+\frac{1}{z-\alpha_{2}}\right)\right) u=0$,
the Riemann hypergeometric equation. As an exercise, show $f_{0}(\lambda)=\lambda(\lambda-$

1) $+a_{0}^{(1)} \lambda+a_{0}^{(2)}$.

## 7 Hypergeometric equation

### 7.1 Riemann hypergeometric equation

A Riemann hypergeometric equation

$$
u^{\prime \prime}+\left(\frac{A_{1}}{z-\alpha_{1}}+\frac{A_{3}}{z-\alpha_{3}}\right) u^{\prime}+\left(\frac{B_{1}}{\left(z-\alpha_{1}\right)^{2}}+\frac{B_{2}}{\left(z-\alpha_{2}\right)^{2}}+C_{1}\left(\frac{1}{z-\alpha_{1}}+\frac{1}{z-\alpha_{2}}\right)\right) u=0,
$$

has indicial equations are

$$
\begin{aligned}
& z=\alpha_{1}: \lambda^{2}+\left(A_{1}-1\right) \lambda+B_{1}=0 \\
& z=\alpha_{2}: \lambda^{2}+\left(A_{2}-1\right) \lambda+B_{2}=0 \\
& z=\infty: \lambda^{2}+\left(1-A_{1}-A_{2}\right) \lambda+B_{1}+B_{2}+\left(\alpha_{1}-\alpha_{2}\right) C_{1}=c
\end{aligned}
$$

$A_{1}, A_{2}, B_{2}, B_{2}, C_{1}$ determine the exponents $\lambda_{ \pm}^{j}$. Conversely,

$$
\begin{aligned}
& A_{j}=1-\lambda_{+}^{(j)}-\lambda_{-}^{(j)} \\
& B_{j}=\lambda_{+}^{(j)} \lambda_{-}^{(j)} \\
& C_{j}=\frac{\lambda_{+}^{(\infty)} \lambda_{-}^{(\infty)}-\left(B_{1}+B_{2}\right)}{\alpha_{1}-\alpha_{2}}
\end{aligned}
$$

Theorem 7.1. The Riemannian equation is completely determined by the seven parameters $\alpha_{1}, \alpha_{2}$ and the exponents (note that the exponents sum up to 1), which is typically presented as

$$
\left(\begin{array}{ccc}
\alpha_{1} & \alpha_{2} & \infty \\
\lambda_{+}^{(1)} & \lambda_{+}^{(2)} & \lambda_{+}^{(\infty)} \\
\lambda_{-}^{(1)} & \lambda_{-}^{(2)} & \lambda_{-}^{(\infty)}
\end{array}\right)
$$

and is called a Riemann scheme. The collection of all solutions to the Riemann equations is called the Riemann symbol, denoted by

$$
P\left(\begin{array}{ccc}
\alpha_{1} & \alpha_{2} & \infty \\
\lambda_{+}^{(1)} & \lambda_{+}^{(2)} & \lambda_{+}^{(\infty)} ; z \\
\lambda_{-}^{(1)} & \lambda_{-}^{(2)} & \lambda_{-}^{(\infty)}
\end{array}\right)
$$

### 7.1.1 Möbius transformation

Recall that a Möbius transformation is an automorphism of $\overline{\mathbb{C}}$ of the form

$$
z \mapsto \frac{a z+b}{c z+d}
$$

where $a, b, c, d \in \mathbb{C}, a d-b c \neq 0$. This is an action by $\mathrm{GL}(n, \mathbb{C})$ with kernel the scalar matrices so is an action by $\operatorname{PSL}(2, \mathbb{C})$.

Theorem 7.2. Characteristic exponents are invariant under Möbius tran-
formation, i.e.

$$
P\left(\begin{array}{lll}
\alpha_{1} & \alpha_{2} & \infty \\
\lambda_{+}^{(1)} & \lambda_{+}^{(2)} & \lambda_{+}^{(\infty)}
\end{array} ; z\right)=P\left(\begin{array}{ccc}
h\left(\alpha_{1}\right) & h\left(\alpha_{2}\right) & h(\infty) \\
\lambda_{-}^{(1)} & \lambda_{-}^{(2)} & \lambda_{-}^{(\infty)}
\end{array}\right)=\left(\begin{array}{cc}
(1) & \lambda_{+}^{(2)} \\
\lambda_{+}^{(\infty)} & ; t=h(z) \\
\lambda_{-}^{(1)} & \lambda_{-}^{(2)}
\end{array} \lambda_{-}^{((\infty)}\right) .
$$

Proof. For simplicity we prove the result for $t=h(z)=\frac{z-\alpha_{1}}{\alpha_{2}-\alpha_{1}}$ so $z=\alpha_{1}, \alpha_{2}, \infty$ corresponds to $t=0,1, \infty$.

$$
\frac{d}{d z}=\frac{1}{\alpha_{2}-\alpha_{1}} \frac{d}{d t}
$$

so

$$
\frac{d^{2} u}{d t^{2}}+\left(\frac{A_{1}}{t}+\frac{A_{2}}{t-1}\right) \frac{d u}{d t}+\left(\frac{B_{1}}{t^{2}}+\frac{B_{2}}{(t-1)^{2}}+\left(\alpha_{2}-\alpha_{1}\right) C_{1}\left(\frac{1}{t}-\frac{1}{t-1}\right)\right) u=0 .
$$

The indicial equations at $t=0,1$ are the same so $\lambda_{ \pm}^{(1)}, \lambda_{ \pm}^{(2)}$ are the same. For $t=\infty$, note that

$$
(0-1)\left(\alpha_{2}-\alpha_{1}\right)=\alpha_{1}-\alpha_{2}
$$

so again it is the same.

Corollary 7.3. We can always reduce a Riemann hypergeometric equation to the form

$$
P\left(\begin{array}{ccc}
0 & 1 & \infty \\
\lambda_{+}^{(0)} & \lambda_{+}^{(1)} & \lambda_{-}^{(\infty)}
\end{array} ; z\right) .
$$

### 7.1.2 Gauge transformation

Consder the gauge transformation $u(z)=z^{p}(1-z)^{q} v(z), p, q \in \mathbb{C}$ that changes the exponent:

$$
\begin{gathered}
u(z)=z^{\lambda_{ \pm}^{(0)}} \operatorname{reg}(z) \mapsto v(z)=z^{\lambda_{ \pm}^{(0)}-p} \operatorname{reg}(z) \\
u(z)=(1-z)^{\lambda_{ \pm}^{(1)}} \operatorname{reg}(z-1) \mapsto v(z)=(1-z)_{ \pm}^{\lambda_{ \pm}^{(1)}-q} \operatorname{reg}(z-1) \\
u(z)=z^{-\lambda_{ \pm}^{(\infty)}} \operatorname{reg}\left(\frac{1}{z}\right) \mapsto v(z)=z^{-\lambda_{ \pm}^{(\infty)}-p-q} \operatorname{reg}\left(\frac{1}{z}\right)
\end{gathered}
$$

so

Thus we can always reduce to the form

$$
P\left(\begin{array}{ccc}
0 & 1 & \infty \\
0 & 0 & \alpha ; z \\
1-\gamma \gamma-\alpha-\beta & \beta
\end{array}\right)
$$

whose correspondig equation is

$$
u^{\prime \prime}+\left(\frac{\gamma}{z}+\frac{\alpha+\beta-\beta+1}{z-1}\right) u^{\prime}+\alpha \beta\left(\frac{1}{z-1}-\frac{1}{z}\right) u=0
$$

i.e.

$$
z(1-z) u^{\prime \prime}+(\gamma-(\alpha+\beta+1) z) u^{\prime}-\alpha \beta u=0 .
$$

This is called the Gauss hypergeometric equation.

### 7.2 Gauss hypergeometric equation

### 7.2.1 Local representation of solutions

We look for a series solution first at $z=0 . \lambda_{+}^{(0)}=0$ so we expect a Taylor series $u(z)=\sum_{n=0}^{\infty} c_{n} z^{n}$. Substitute to get the recurrence relation

$$
c_{n+1}=\frac{(m+\alpha)(m+\beta)}{(n+1)(n+\gamma)} c_{n}
$$

where we require $\gamma \neq 0,-1,-2, \ldots$ so that Frobenius method works. Then

$$
c_{n}=\frac{(\alpha)_{n}(\beta)_{n}}{n!(\gamma)_{n}} c_{0}
$$

where the subscript $n$ is the Pochhammer symbol and is defined as

$$
(\alpha)_{0}=1,(\alpha)_{n}=\alpha(\alpha+1) \cdots(\alpha+n-1)
$$

Note that $(\alpha)_{n}=0$ for $\alpha=0,-1, \ldots,-n+1$. Set $c_{0}=1$, we get

$$
u_{1}^{(0)}(z)={ }_{2} F_{1}(\alpha, \beta, \gamma ; z)=\sum_{n=0}^{\infty} \frac{(\alpha)_{n}(\beta)_{n}}{(\gamma)_{n} n!} z^{n}
$$

convergent for $|z|<1$, called the hypergeometric series.
To get the second solution: use gauge $u=z^{1-\gamma} v$ so

$$
P\left(\begin{array}{ccc}
0 & 1 & \infty \\
0 & 0 & \alpha \\
1-\gamma & \gamma-\alpha-\beta & \beta
\end{array}\right)=z^{1-\gamma} P\left(\begin{array}{ccc}
0 & 1 & \infty \\
0 & 0 & \alpha-\gamma+1 ; z \\
\gamma-1 & \gamma-\alpha-\beta & \beta-\gamma+1
\end{array}\right)
$$

so the other solution is

$$
u_{2}^{(0)}(z)=z^{1-\gamma}{ }_{2} F_{1}(\alpha-\gamma+1, \beta-\gamma+1,2-\gamma ; z)
$$

where $\gamma \neq 2,3, \ldots$.
To get the solution at 1 , apply the Möbius transformation $t=1-z$ so

$$
P\left(\begin{array}{ccc}
0 & 1 & \infty \\
0 & 0 & \alpha ; z \\
1-\gamma \gamma-\alpha-\beta & \beta
\end{array}\right)=P\left(\begin{array}{ccc}
0 & 1 & \infty \\
0 & 0 & \alpha \\
\gamma-\alpha-\beta & 1-\gamma & \beta
\end{array}\right)
$$

so by a change of variable $\gamma-\alpha-\beta=1-\gamma_{1}, \alpha_{1}=\alpha, \beta_{1}=\beta$, from the solutions at $z=0$ we get
$u_{1}^{(1)}(z)=F(\alpha, \beta, \alpha+\beta-\gamma+1 ; 1-z) \quad \gamma-\alpha-\beta \neq 1,2, \ldots,|1-z|<1$
$u_{1}^{(2)}(z)=(1-z)^{\gamma-\alpha-\beta} ;{ }_{2} F_{1}(\gamma-\beta, \gamma-\alpha, 1+\gamma-\alpha-\beta ; 1-z) \quad a \gamma-\alpha-\beta \neq-1,-2, \ldots$
To get the solution at $\infty$ we use $z=\frac{1}{t}$ and after another computation

$$
\begin{aligned}
& u_{1}^{(\infty)}(z)=z^{-\alpha}{ }_{2} F_{1}\left(\alpha, 1+\alpha-\gamma, 1+\alpha-\beta ; \frac{1}{z}\right) \quad \alpha-\beta \neq-1,-2, \ldots,|z|>1 \\
& u_{1}^{(\infty)}(z)=z^{-\beta}{ }_{2} F_{1}\left(\beta, 1+\beta-\gamma, 1-\alpha+\beta ; \frac{1}{z}\right) \quad \alpha-\beta \neq 1,2, \ldots
\end{aligned}
$$

### 7.2.2 Integral representation

Recall the gamma function is defined by the Euler integral

$$
\Gamma(p)=\int_{0}^{\infty} t^{p-1} e^{-t} d t
$$

for $\operatorname{Re} p>0$. Also Beta function

$$
B(p, q)=\int_{0}^{1} t^{p-1}(1-t)^{q-1} d t
$$

for $\operatorname{Re} p>0, \operatorname{Re} q>0, \arg t=0, \arg (1-t)=0$ which satisfies

$$
B(p, q)=\frac{\Gamma(p) \Gamma(q)}{\Gamma(p+q)}
$$

Proposition 7.4. Suppose $\operatorname{Re} \gamma>\operatorname{Re} \beta>0,|z|<1$, we have the following integral representation of hypergeometric function:

$$
{ }_{2} F_{1}(\alpha, \beta, \gamma ; z)=\frac{\Gamma(\gamma)}{\Gamma(\beta) \Gamma(\gamma-\beta)} \int_{0}^{1} t^{\beta-1}(1-t)^{\gamma-\beta-1}(1-z t)^{-\alpha} d t
$$

where $\arg t=0, \arg (1-t)=0,|\arg (1-z t)|<\frac{\pi}{2}$ for $0<t<1$. This is the Euler representation of hypergeometric functions.

## Proof.

$$
\begin{aligned}
{ }_{2} F_{1}(\alpha, \beta, \gamma ; z) & =\sum_{n=0}^{\infty} \frac{(\beta)_{n}}{(\gamma)_{n}} \frac{(\alpha)_{n} z^{n}}{n!} \\
& =\sum_{n} \underbrace{\Gamma(\beta) \Gamma(\gamma+n)}_{\frac{\Gamma(\gamma)}{\Gamma(\beta)} \frac{\Gamma(n+\beta, \gamma-\beta)}{\Gamma(\gamma-\beta)}} \frac{\Gamma(\beta+n)}{n!} \\
& =\frac{\Gamma(\gamma)}{\Gamma(\beta) \Gamma(\gamma-\beta)} \sum_{n} \int_{0}^{1} t^{\beta+m-1}(1-t)^{\gamma-\beta-1}\binom{-\alpha}{n}(-z)^{n} d t \\
& =\frac{\Gamma(\gamma)}{\Gamma(\beta) \Gamma(\gamma-\beta)} \int_{0}^{1} t^{\beta-1}(1-t)^{\gamma-\beta-1} \underbrace{\sum_{n}\binom{-\alpha}{n}(-t z)^{n}}_{=(1-t z)^{-\alpha}} d t \quad \text { as }|z|<1
\end{aligned}
$$

where the expansion for $(1-t z)^{-\alpha}$ is valid for $0<z t<1$ and $z$ real. ...
In the Euler representation of ${ }_{1} F_{2}$, the integrand has a branch point at $z=\frac{1}{t}$ for $0 \leq t \leq 1$. In the $z$-plane, as $t$ varies, the locus of $\frac{1}{t}$ is the real axis greater than 1 . Thus we can analytically continue ${ }_{2} F_{1}$ to $\mathbb{C} \backslash[1, \infty)$.

### 7.2.3 Solution to connection problem

Suppose $\operatorname{Re}(\gamma-\alpha-\beta)>0, \operatorname{Re} \gamma>\operatorname{Re} \beta>0$. Then

$$
\begin{aligned}
{ }_{2} F_{1}(\alpha, \beta, \gamma ; 1) & =\frac{\Gamma(\gamma)}{\Gamma(\beta) \Gamma(\gamma-\beta)} \int_{0}^{1} t^{\beta-1}(1-t)^{\gamma-\alpha-\beta-1} d t \\
& =\frac{\Gamma(\gamma)}{\Gamma(\beta) \Gamma(\gamma-\beta)} B(\beta, \gamma-\alpha-\beta) \\
& =\frac{\Gamma(\gamma)}{\Gamma(\beta) \Gamma(\gamma-\beta)} \frac{\Gamma(\beta) \Gamma(\gamma-\alpha-\beta)}{\Gamma(\beta+\gamma-\alpha-\beta)} \\
& =\frac{\Gamma(\gamma) \Gamma(\gamma-\alpha-\beta)}{\Gamma(\gamma-\alpha) \Gamma(\gamma-\beta)}
\end{aligned}
$$

This is called the Gauss-Kummer formula. This formula holds by analytic continuation with $\gamma \neq 0,-1,-2, \ldots$ and $\gamma-\alpha-\beta \neq 0,-1, \ldots$.

Recall that we have the fundamental solution to the Gauss hypergeometric equation. We write them in row matrices $\left[u_{1}^{(0)}, u_{2}^{(0)}\right]=Y^{(0)}$ etc. The main problem is to compute the connection matrices $C_{01}$ and $C_{0 \infty}$ where $Y^{(0)}=$ $Y^{(1)} C_{01}, Y^{(0)}=Y^{(\infty)} C_{0 \infty}$. We choose branch cuts so $|\arg z|<\pi,|\arg (1-z)|<$ $\pi$. As an example computation, suppose $C_{01}=\left(\begin{array}{cc}a & c \\ b & d\end{array}\right)$. Then

$$
u_{1}^{(0)}(z)=a u_{1}^{(1)}(z)+b u_{2}^{(1)}(z)
$$

Substitute $z=1$ into the equation and we get

$$
{ }_{2} F_{1}(\alpha, \beta, \gamma ; 1)=a \cdot \underbrace{{ }_{2} F_{1}(\alpha, \beta, \alpha+\beta-\gamma+1 ; 0)}_{=1}+b \cdot 0
$$

where for the last term to be 0 we require $\operatorname{Re}(\gamma-\alpha-\beta) \geq 0$. Apply GaussKummer we can find $a$. By substituting $z=0$ we get

$$
1=a \cdot{ }_{2} F_{1}(\cdots ; 1)+b(1-0)^{\gamma-\alpha-\beta}{ }_{2} F_{1}(\cdots ; 1) .
$$

Note as $z \rightarrow 1,1-z \rightarrow 0$ so $\arg (1-z) \rightarrow 0$ so the exponential term goes to 1 . In this way we can solve for $b$.

This illustrates the general strategy of solving the connection problem: we find an integral representation for the solutions, analytically continue them and evaluate them in some "common point" $z_{0}$.

### 7.2.4 Monodromy

Since the connection matrices are known, we can compute monodromy using from local representations: under $z \mapsto z e^{2 \pi 1}$, recall that $u_{1}^{(0)}$ has no monodromy and $u_{2}^{(0)}$ has the exponential term $z^{1-\gamma}$, we have

$$
Y^{(0)} \mapsto Y^{(0)}\left(\begin{array}{ll}
1 & \\
& e^{-2 \pi i \gamma}
\end{array}\right)
$$

Similarly

$$
\begin{aligned}
& Y^{(1)} \mapsto Y^{(1)}\left(\begin{array}{cc}
1 & \\
& e^{2 \pi i(\gamma-\beta-a)}
\end{array}\right) \\
& Y^{(\infty)} \mapsto Y^{(\infty)}\left(\begin{array}{cc}
e^{2 \pi i \alpha} & \\
& e^{2 \pi i \beta}
\end{array}\right)
\end{aligned}
$$

Denote the local monodromy matrices by $\tilde{M}_{0}$ etc, we have that for $Y^{(0)}$ the monodromy matrices are

$$
M_{0}=\tilde{M}_{0}, M_{1}=C_{01}^{-1} \tilde{M}_{1} C_{01}, M_{\infty}=C_{0 \infty}^{-1} \tilde{M}_{\infty} C_{0 \infty}
$$

We present here another approach that works for Gauss hypergeometric equations. Choose as fundamental system $u_{1}^{(0)}, u_{1}^{(1)}$ (they are linearly independent). Then immediately we know the monodromy matrices have the form

$$
M_{0}=\left(\begin{array}{ll}
1 & * \\
0 & *
\end{array}\right), M_{1}=\left(\begin{array}{cc}
* & 0 \\
* & 1
\end{array}\right)
$$

Recall that the exponents are $1-\gamma$ at $z=0$ and $\gamma-\alpha-\beta$ at $z=1$, we have

$$
M_{0}=\left(\begin{array}{cc}
1 & x \\
0 & e^{2 \pi i(1-\gamma)}
\end{array}\right), M_{1}=\left(\begin{array}{cc}
e^{2 \pi i(\gamma-\alpha-\beta)} & 0 \\
y & 1
\end{array}\right)
$$

for some $x$ and $y$. As $M_{1} M_{0}=M_{\infty}^{-1}$ and $M_{\infty}^{-1}$ has eigenvalues $e^{-2 \pi i \alpha}, e^{-2 \pi i \beta}$, taking trace we can find $x y$ (in terms of $\alpha, \beta$ and $\gamma$ ). Note that this is all we can extract: if we change the fundamental system by a linear transformation to $\left[r u_{0}^{(0)}, s u^{(u)}\right]$ then the monodromy matrix $M_{j}$ transforms to

$$
\left(\begin{array}{ll}
r & \\
& s
\end{array}\right)^{-1} M_{j}\left(\begin{array}{ll}
r & \\
& s
\end{array}\right)=\left(\begin{array}{cc}
* & * \frac{r}{s} \\
* \frac{s}{r} & *
\end{array}\right)
$$

Thus the invariant is $x y$ and $x, y$ individually have freedom (by diagonal conjugation). We conclude by the following "rigidity" property of the Gauss hypergeometric equation: the equation is equivalent to the characteristic exponents, and the characteristic exponent determines monodromy (up to diagonal conjugation) and vice versa. Thus in turn the equation and the monodoromy determine each other.

$$
\text { equation } \leftrightarrow \text { characteristic exponent } \leftrightarrow \text { monodromy }
$$

### 7.3 Some further results

Companion system Let $y=\left(y_{1}, y_{2}\right)^{T}$ where $y_{1}=u, y_{2}=(z-1) u^{\prime}$. Then $y$ satisfies the following system

$$
\frac{d y}{d z}=(\frac{1}{z} \underbrace{\left(\begin{array}{cc}
0 & 0 \\
-\alpha \beta & -\gamma+
\end{array}\right)}_{A_{0}}+\frac{1}{z-1} \underbrace{\left(\begin{array}{cc}
0 & 1 \\
0 & \gamma-\alpha-\beta
\end{array}\right)}_{A_{1}}) y
$$

The eigenvalues of $A_{0}$ are $0,-\gamma$, the eigenvalues of $A_{1}$ are $0, \gamma-\alpha-\gamma$ and the eigenvalues of $A_{\infty}=-A_{0}-A_{1}$ are $\alpha, \beta$

## Monodromy

Theorem 7.5. Any irreducible representation

$$
\pi_{1}(\overline{\mathbb{C}} \backslash\{0,1, \infty\}) \rightarrow \operatorname{GL}(2, \mathbb{C})
$$

is realised by the Guass equation.

Theorem 7.6. Any irreducible represdentation is realised by the monodromy of

$$
\frac{d y}{d z}=\left(\frac{1}{z}\left(\begin{array}{ll}
* & * \\
0 & *
\end{array}\right)+\frac{1}{z-1}\left(\begin{array}{ll}
* & * \\
0 & *
\end{array}\right)\right) y .
$$

Theorem 7.7. Any $2 \times 2$ Fuchsian system

$$
\frac{d y}{d z}=\left(\frac{A_{0}}{z}+\frac{A_{1}}{z-1}\right) y
$$

with an irreducible monodromy is meromophically equivalent to the Gauss equation.

## 8 Poincaré asymptotics

As a motivating example, consider the system, which is confusingly also called the Euler equation,

$$
y^{\prime}+\frac{1}{z^{2}} y=\frac{1}{z}
$$

Note that $z=0$ is non-Fuchsian. It has homogeneous solution $y_{h}(z)=c e^{1 / z}$. By variation of parameters, subsitute $y_{p}(z)=c(z) e^{1 / z}$ into the equation and solve to get

$$
c(z)=\int_{z_{0}}^{z} \frac{e^{-1 / s}}{s} d s
$$

so by a change of variable

$$
y_{p}(z)=e^{1 / z} \int_{1 / z}^{1 / z_{0}} \frac{e^{-t}}{t} d t
$$

We can take $z>z_{0}>0$. If we take the limit $z_{0} \rightarrow 0^{+}$, we get

$$
y(z)=e^{1 / z} \int_{1 / z}^{\infty} \frac{e^{-t}}{t} d t
$$

which is real for $z>0$ and can be analytically continued to $\operatorname{Re} z>0$. We call the integral

$$
\operatorname{Ei}(z)=\int_{z}^{\infty} \frac{e^{-t}}{t} d t
$$

the exponential integral. The solution $y(z)=e^{1 / 2} \operatorname{Ei}(z)$ is called Euler function and is denote by $E(z)$.

We look for a formal solution $y_{f}(z)=\sum_{m=0}^{\infty} c_{m} z^{m}$. Substitute into the equation and solve to get

$$
y_{f}(z)=\sum_{m}(-1)^{m} m!z^{m+1}
$$

which is divergent. But $E(z) \sim y_{f}(z)$ when $z \rightarrow 0$ in $-\frac{3 \pi}{2}<\arg z<\frac{3 \pi}{2}$, meaning that $\left|\sum_{m=0}^{N} c_{m} z^{m}-E(z)\right|$ is

- small for $|z|$ small and $N$ fixed, and
- small for large $N$ and $|z|$ fixed.

Poincaré posed the following problem. Take as example two series

$$
\sum_{n=0}^{\infty} \frac{1000^{n}}{n!}, \sum_{n=0}^{\infty} \frac{n!}{1000^{n}}
$$

The first is convergence while the second is convergent. The question is: suppose we want to compute the partial sum of the first, say, 1000 terms. However terms do we need to sum so that the error is within $10^{-6}$ ? In other words for what $N \leq 1000$ we have

$$
\sum_{n=0}^{1000} s_{n}-\sum_{n=0}^{N} s_{n}<10^{-6} ?
$$

With the help of Stirling formula

$$
m!\sim m^{m} \sqrt{2 \pi m} e^{-m},
$$

for the first series we need $N=1000$ but for the second we need $N=3$. Thus Poincaré said that the second series, although divergent for mathematicians, is "convergent" for astronomers.

Let $s(\alpha, \beta)=\{z \in \widetilde{\mathbb{C} \backslash\{0\}}: \alpha<\arg z<\beta\}$ be an open sector (with vertex $z=0$ or $z=\infty)$. We can also define analogously a closed sector $\bar{s}(\alpha, \beta)$ by using nonstrict inequalities.

Definition (asymptotic series). Let $S$ be a sector with vertex at $z=0$. Let $f: S \rightarrow \mathbb{C}$ be a function and $\sum_{k=0}^{\infty} a_{k} z^{k}$ be a formal series. We say $f(z)$ is asymptotic to the formal series, written $f(z) \sim \sum_{k=0}^{\infty} a_{k} z^{k}$, for $z \rightarrow 0$ in $S$ if for every closed subsector $S^{\prime} \subseteq S$ and for every $m \in \mathbb{N}$, exists $C\left(m, S^{\prime}\right)>0$ such that

$$
\left|f(z)-\sum_{k=0}^{m} a_{k} z^{k}\right| \leq C\left(m, S^{\prime}\right)|z|^{m+1}
$$

for all $z \in S^{\prime}$.
We say the asymptotic series is uniform in $S$ if $C\left(m, S^{\prime}\right)$ does not depend on $S^{\prime}$.

Of course if $S$ is closed then we only need to check $S^{\prime}=S$. The key point is that if the sector is open then $C\left(m, S^{\prime}\right)$ may depend on $S^{\prime}$ as $S^{\prime}$ "approaches" $S$.
Exercise. Verify that the definition above is equivalent to saying that for every closed $S^{\prime} \subseteq S$,

$$
\frac{1}{z^{m}}\left(f(z)-\sum_{k=0}^{m} a_{k} z^{k}\right) \rightarrow 0
$$

as $z \rightarrow 0$ in $S^{\prime}$.
At $z=\infty$, we say $f(z) \sim \sum_{k=0}^{\infty} a_{k} z^{-k}$ as $z \rightarrow \infty$ in $S$ if for every closed $S^{\prime} \subseteq S$ and every $m \in \mathbb{N}$, exists $C\left(m, S^{\prime}\right)>0$ such that

$$
\left|f(z)-\sum_{k=0}^{m} \frac{a_{k}}{z^{k}}\right| \leq \frac{C\left(m, S^{\prime}\right)}{|z|^{m+1}}
$$

for $z \in S^{\prime}$.
Some properties of asymptotic series:

1. if $f$ has asymptotic expansion then it is unique. This is guaranteed by the uniqueness of limit as $a_{0}=\lim _{z \rightarrow 0, z \in S^{\prime}}=a_{0}$ and inductively all $a_{m}$ 's are uniquely determined.
2. not all functions have asymptotic expansion in a given sector. For example consider $f(z)=e^{1 / z}$ and consider for example the sector $S\left(-\frac{\pi}{4}, \frac{\pi}{4}\right)$. Then $f(z) \rightarrow \infty$ as $z \rightarrow 0$.
3. $\sum a_{k} z^{k}$ can be (is) asymptotic expansion of infinitely many functions. For example the zero formal series is asymptotic to $f(z)=e^{-\frac{1}{z^{\sigma}}} \sim 0$ in $-\frac{\pi}{2 \sigma}<\arg z<\frac{\pi}{2 \sigma}$ for all $\sigma>0$, as $\lim \frac{f(z)}{z^{m}}=0$ for all $m$ in the sector.

Proposition 8.1. Suppose $f(z) \sim \sum_{k=0}^{\infty} a_{k} z^{k}$ in S. If the opening angle is $>2 \pi$ and $f$ is single-valued at $z=0$ then the series converges.

Proof. We know $a_{0}=\lim _{z \rightarrow 0, z \in S} f(z)$. The limit is along any direction in $\mathbb{C} \backslash\{0\}$, so $z=0$ is a removable singularity. Thus $f$ is holomorphic at $z=0$ so has a convergent Taylor expnansion which necessarily coincide with the given series.

Wasow: Asymptotic expansions for ODE
Algebraic properties

1. Let $f(z) \sim \sum_{k=0}^{\infty} a_{k} z^{k}, g(z) \sim \sum_{k=0}^{\infty} b_{k} z^{k}$ as $z \rightarrow 0$ in $S$. Let $\alpha, \beta \in \mathbb{C}$, then

$$
\begin{aligned}
\alpha f(z)+\beta g(z) & \sim \sum_{k=0}^{\infty}\left(\alpha a_{k}+\beta b_{k}\right) z^{k} \\
f(z) \cdot g(z) & \sim \sum_{\ell=0}^{\infty}\left(\sum_{k=0}^{\ell} a_{k} b_{\ell-k}\right) z^{\ell} \\
\frac{1}{f(z)} & \sim \sum_{k=0}^{\infty} c_{k} z^{k} \text { if } a_{0} \neq 0, \text { where } \sum a_{k} z^{k} \cdot \sum c_{k} z^{k}=1
\end{aligned}
$$

2. Suppose $f \sim \sum a_{k} z^{k}$ in $S_{f}, g(z)=\sum b_{k} z^{k}$ in $S_{g}$. Let $\tilde{f}(z)=f(z)-a_{0} \sim$ $\sum_{k=1}^{\infty} a_{k} z^{k}$. Assume $\tilde{f}\left(S_{f}\right) \subseteq S_{g}$. Then

$$
g(\tilde{f}(z)) \sim \sum_{h=0}^{\infty} b_{h}\left(\sum_{k=1}^{\infty} a_{k} z^{k}\right)^{h}
$$

in $S_{f}$.
Analytic properties:

1. Suppose $f$ is holomorphic in $S$ with centre $z=0$ and $f(z) \sim \sum a_{k} z^{k}$ in $S$. Then

$$
\int_{0}^{z} f(\zeta) d \zeta \sim \sum a_{k} \int_{0}^{z} \zeta^{k} d \zeta=\sum \frac{a_{k}}{k+1} z^{k+1}
$$

in $S$.
If $f(z) \sim \sum \frac{a_{k}}{z^{k}}$ as $z \rightarrow \infty$ then

$$
\int_{\infty}^{z}\left(f(\zeta)-a_{0}-\frac{a_{1}}{\zeta} d \zeta\right) \sim \sum_{k=2}^{\infty} a_{k} \int_{\infty}^{z} \frac{d \zeta}{\zeta^{k}}=\sum_{k=2}^{\infty} \frac{a_{k}}{1-k} z^{1-k}
$$

as $z \rightarrow \infty$ in $S$.
2. Suppose $f(z)$ is holomorphic on $S$ at $z=0$ and $f(z) \sim \sum a_{k} z^{k}$ in $S$, a sector with interior (i.e. not a ray). Then

$$
\frac{d f(z)}{d z} \sim k a_{k} z^{k-1}
$$

in every closed subsector $S^{\prime} \subseteq S$.
3. Note that if we add the condition that $f(z)$ has a uniform asymptotic series, it is not true that we get a uniform asymptotic for $\frac{d f(z)}{d z}$.
4. Suppose $f(z) \sim \sum a_{k} z^{k}$ as $z \rightarrow 0$ in $S$, a sector with interior. Then the limit $\lim _{z \rightarrow 0} f^{(k)}(z)$ exists and equals to $k!a_{k}$ in every $S^{\prime} \subseteq S$.
5. Suppose $f(z)$ is holomorphic in $S$ and assume $\lim _{z \rightarrow 0} f^{(k)}(z)=f_{k}$. Then

$$
f(z) \sim \sum \frac{f_{k}}{k!} z^{k}
$$

as $z \in 0$ in $S$.

Theorem 8.2 (Borel-Ritt). For every formal series $\sum_{k=0}^{\infty} a_{k} z^{k}$ and every sector $S$ (open or closed) at $z=0$, there exists $f(z)$ holomorphic in $S$ such that $f(z) \sim \sum a_{k} z^{k}$ in $S$.
example of computation of asymptotic expansion Recall the exponential integral

$$
\operatorname{Ei}(z)=\int_{z}^{\infty} \frac{e^{-t}}{t} d t, z>0
$$

Proposition 8.3. The exponential integral defines an analytic function in $\{z \in \mathbb{C} \backslash\{0\}:-\pi<\arg z<\pi\}$. It is real for $z>0$ and

$$
\operatorname{Ei}(z)=-\ln z-\gamma-\sum_{n=1}^{\infty} \frac{(-1)^{n+1}}{n \cdot n!} z^{n}
$$

where $\ln z$ is the principal branch, i.e. $\ln z \in \mathbb{R}$ for $z>0, \gamma$ is the EulerMascheroni constant, and the series converges uniformly in every compact subset of $\mathbb{C}$.

Proof.

$$
\begin{aligned}
\operatorname{Ei}(z) & =\int_{1}^{z} \frac{e^{-t}}{t} d t+\int_{z}^{1} \frac{e^{-t}}{t} d t \\
& =c_{1}+\int_{z}^{1} \sum_{0}^{\infty} \frac{(-1)^{n} t^{n-1}}{n!} d t \\
& =c_{1}+\int_{z}^{1} \frac{d t}{t}+\sum_{n=1}^{\infty} \frac{(-1)^{n}}{n!} \int_{z}^{1} t^{n-1} d t \\
& =c_{2}-\ln z+\sum_{n=1}^{\infty} \frac{(-1)^{n+1}}{n!} \frac{z^{n}}{n}
\end{aligned}
$$

so $\operatorname{Ei}(z)+\ln z$ is analytic at $z=0$. As $\operatorname{Ei}(z) \in \mathbb{R}$ for $z>0, \ln z$ takes the principal branch.

To find $c_{2}$, integrate by parts

$$
\operatorname{Ei}(z)=-e^{-z} \ln z+\int_{z}^{\infty} e^{-t} \ln d t
$$

so

$$
c_{2}=\lim _{z \rightarrow 0}(\operatorname{Ei}(z)+\ln z)=\lim _{z \rightarrow 0}\left[\left(1-e^{-z} \ln z\right)+\int_{z}^{\infty} e^{-t} \ln t d t\right]=\int_{0}^{\infty} e^{-t} \ln t
$$

which we call $-\gamma$.
Exercise. Show

$$
-\gamma=\left.\frac{d \ln \Gamma(z)}{d z}\right|_{z=1},-\gamma=\int_{0}^{\infty} e^{-t} \ln t d t
$$

From the expression we also know that the analytic continuation of $\operatorname{Ei}(z)$ on $\widetilde{\mathbb{C} \backslash\{0\}}$ has monodromy

$$
\operatorname{Ei}\left(z e^{2 \pi i}\right)=\operatorname{Ei}(z)-2 \pi i .
$$

Proposition 8.4. $\operatorname{Ei}(z)$ defined on $\widetilde{\mathbb{C} \backslash\{0\}}$ has the following asymptotic representation:

$$
z e^{z} \operatorname{Ei}(z) \sim \sum_{n=0}^{\infty}(-1)^{n} n!z^{-n}
$$

as $z \rightarrow \infty$ in $S=\left\{z \in \widetilde{\mathbb{C} \backslash\{0\}}:-\frac{3 \pi}{2}<\arg z<\frac{3 \pi}{2}\right.$.
Proof.

$$
\begin{aligned}
\int_{z}^{\infty} \frac{e^{-t}}{t} d t & =\frac{e^{-z}}{z}-\int_{z}^{\infty} \frac{e^{-t}}{t^{2}} d t \\
& =\frac{e^{-z}}{z}-\frac{e^{-z}}{z^{2}}+2 \int_{z}^{\infty} \frac{e^{-t}}{t^{3}} d t \\
& =\cdots \\
& =\frac{e^{-z}}{z}[\sum_{n=0}^{m} \frac{(-1)^{n} n!}{z^{n}}+\underbrace{(-1)^{m+1}(m+1)!z \int_{z}^{\infty} \frac{e^{z-t}}{t^{m+2}} d t}_{R_{m}(z)}]
\end{aligned}
$$

so

$$
z e^{z} \operatorname{Ei}(z)=\sum_{n=0}^{m} \frac{(-1)^{n} n!}{z^{n}}+R_{m}(z)
$$

Check $\left|R_{m}(z)\right| \leq C(m)|z|^{m+1}$ :

$$
\begin{aligned}
\left|R_{m}(z)\right| & =(m+1)!|z|\left|\int_{z}^{\infty} \frac{e^{z-t}}{t^{m+2}} d t\right| \\
& =(m+1)!|z|\left|\frac{1}{z^{m+2}}-(m+2) \int_{z}^{\infty} \frac{e^{z-t}}{t^{m+3}} d t\right| \\
& \leq(m+1)!|z|^{-m-1}+(m+2)!\left|\int_{z}^{\infty} \frac{e^{z-t}}{t^{m+3}} d t\right|
\end{aligned}
$$

For $z>0$, we have estimate

$$
\left|\int_{z}^{\infty} \frac{e^{z-t}}{t^{m+3}} d t\right| \leq \int_{z}^{\infty} \frac{d t}{t^{m+3}}=\frac{z^{-m-2}}{m+2}
$$

so

$$
\left|R_{m}(z)\right| \leq 2(m+1)!|z|^{-m-1}
$$

as desired.
Exercise. Show the asymptotic series is valid on $-\pi \leq \arg g \leq \pi$, by deriving

$$
\left|R_{m}(z)\right| \leq(2(m+1)!+(m+2) \pi)|z|^{-m-1}
$$

extension of asymptotics outside $-\pi \leq \arg z \leq \pi$ We have asymptotic series for

$$
\begin{gathered}
e^{z} \operatorname{Ei}(z),-\pi \leq \arg z \leq \pi \\
e^{z} \operatorname{Ei}\left(z e^{2 \pi i}\right),-3 \pi \leq \arg z \leq-\pi \\
e^{z} \operatorname{Ei}\left(z e^{-2 \pi i}\right), \pi \leq \arg z \leq 3 \pi
\end{gathered}
$$

Note that the domains overlap. Recall

$$
e^{z} \operatorname{Ei}(z)=e^{z} \operatorname{Ei}\left(z e^{ \pm 2 \pi i}\right) \pm 2 \pi i e^{2}
$$

Finally note that $e^{z} \sim 0$ as $z \rightarrow \infty$ in $\frac{\pi}{2}+2 k \pi<\arg z<\frac{3 \pi}{2}+2 k \pi$. Thus

$$
e^{z} \operatorname{Ei}(z)-e^{z} \operatorname{Ei}\left(z e^{ \pm 2 \pi i}\right) \sim 0
$$

in the sector. Set $k=0$ and -1 and apply to the three expansions, we get the sector $-\frac{3 \pi}{2}<\arg z<\frac{3 \pi}{2}$.

Note. Note the error term

$$
\left|R_{m}(z)\right| \leq \underbrace{(2(m+1)!+(m+2) \pi)}_{C_{n}}|z|^{-m-1} .
$$

is typical of asymptotic expansions:

1. ..

## $9 \quad$ Singularities of the second kind

We will assume the singularity is at $z=\infty$ so we study the system

$$
\frac{d y}{d z}=z^{r-1} A(z) y
$$

where $A(z)$ is holomorphic for $|z| \geq R$ and $A(z)=\sum \frac{A_{i}}{z^{i}}$ for $|z| \geq R$. If $r=0$ then it is a Fuchsian singularity, if $r \geq 1$ then it is a singularity of the seoncd kind. We assume $A(z) \sim \sum_{i=0}^{\infty} \frac{A_{i}}{z^{i}}$ in $S=S(\alpha, \beta)$ and $|z|>R$. The object we are dealing with in tis chapter is the solution

$$
y(z)=G(z) z^{B} \exp (\Lambda(z))
$$

where $G(z)$ is analytic in a sector at $\infty$ and has an asymptotic expansion in $z^{-1}, \Lambda(z)$ is a polynomial in $z$.

### 9.1 Formal simplification

Suppose $A(z)=\sum_{i=0}^{\infty} \frac{A_{i}}{z^{i}}$ where $A_{0}$ has at least two distinct eigenvalues. Divide the eigenvalues of $A_{0}$ in two sequences $\left(\lambda_{1}, \ldots, \lambda_{p}\right),\left(\lambda_{p+1}, \ldots, \lambda_{m}\right)$ which are disjoint. We can find $G_{0}$ invertible such that $G_{0}^{-1} A_{0} G^{0}$ is block diagonal with each block having the respective sequence of eigenvalues. We look for a gauge transformation $Y(z)=G(z) \tilde{Y}(z), G(z)=I+\sum_{j=1}^{\infty} G_{j} z^{-j}$ such that

$$
\frac{d \tilde{Y}}{d z}=z^{r-1} B(z) \tilde{Y}
$$

which is "simpler".
Substitute into the differential equatation, we get

$$
G^{\prime} \tilde{y}+G \tilde{Y}^{\prime}=z^{r-1} A G \tilde{Y},
$$

i.e.

$$
z^{1-r} G^{\prime}+G B=A G .
$$

Substitute $G=\sum G_{i} z^{-i}, A=\sum A_{i} z^{-i}, B=\sum B_{i} z^{-i}$,

- $\ell=0$ : set $G_{0}=I, B_{0}=A_{0}$,
- $1 \leq \ell \leq r: A-G_{\ell}-G_{\ell} A_{0}=\left(\sum_{j=1}^{\ell-1}\left(G_{j} B_{\ell-j}-A_{\ell-j} G_{j}\right)-A_{\ell}\right)+B_{\ell}$
- $r+1 \leq \ell: A-G_{\ell}-G_{\ell} A_{0}=\left(\sum_{j=1}^{\ell-1}\left(G_{j} B_{\ell-j}-A_{\ell-j} G_{j}\right)-(\ell-r) G_{\ell-r}\right)+B_{\ell}$

Thus the recurrence relation for $G_{\ell}, B_{\ell}$ is

$$
A_{0} G_{\ell}-G_{\ell} A_{0}=K_{\ell}+B_{\ell}
$$

where $K_{\ell}$ is known from previous steps.
Recall

$$
A_{0}=\left(\begin{array}{ll}
* & 0 \\
0 & *
\end{array}\right), B_{\ell}=\left(\begin{array}{cc}
B_{11}^{(\ell)} & B_{12}^{(\ell)} \\
B_{21}^{(\ell)} & B_{22}^{(\ell)}
\end{array}\right)
$$

etc so we get Sylvester equation

$$
A_{i i}^{(0)} G_{i j}^{(\ell)}-G_{i j}^{(\ell)} A_{j j}^{(0)}=K_{i j}^{(\ell)}+B_{i j}^{(\ell)}
$$

For $i=j$, take $B_{i i}^{(\ell)}=-K_{i i}^{(\ell)}$. Then $G_{i i}^{(\ell)}$ is the solution of a homogeneous equation so has a nontrivial solution (we can choose $G_{i i}^{(\ell)}=0$ ). For $i \neq j$, take $B_{i j}^{(\ell)}=0$. As $A_{i i}^{(0)}$ and $A_{j j}^{(0)}$ do not have common eigenvalues, there is a solution $G_{i j}^{(\ell)}$.

In conclusion, we can find gauge transformation $G(z)$ such that $Y(z)=$ $G_{0} G(z) \tilde{Y}(z)$ satisfies

$$
\frac{d \tilde{Y}}{d z}=z^{r-1} B(z) \tilde{Y}(z)
$$

where $B(z)$ is block diagonal.
The block diagonal system admits fundamental matrix solution

$$
\tilde{Y}(z)=\left(\begin{array}{cc}
\tilde{Y}_{1}(z) & 0 \\
0 & \tilde{Y}_{2}(z)
\end{array}\right)
$$

and

$$
\frac{d \tilde{Y}_{i}}{d z}=z^{r-1} B_{i i}(z) \tilde{Y}_{i}
$$

We can repeat the above computation for $\tilde{Y}_{i}$ if $B_{i i}^{(0)}$ has at least two distinct eigenvalues. By repeated applications we conclude that the following holds:

Theorem 9.1. If $A_{0}$ has $s \leq n$ distinct eigenvalues $\lambda_{1}, \ldots, \lambda_{s}$, we can formally determine $G(z)$ such that $Y=G_{0} G(z) \tilde{Y}$ gives

$$
\frac{d \tilde{Y}}{d z}=z^{r-1} B(z) \tilde{Y}
$$

where $B(z)$ is block diagonal with s blocks.
Analytic issue of the solution is addressed by Malmquist (1944) and Sibuya (1962) (see Wasow). We present an easier result.

## Theorem 9.2.

1. Let $A(z)$ be analytic for $|z|>R$ with Taylor series $\sum_{i=0}^{\infty} A_{i} z^{-i}$, absolutely convergent on $|z| \geq \rho>R$.
2. It suffices to suppose that $A(z)$ is analytic in a sector $S$ for opening angle $<\frac{\pi}{r},|z|>R$ with $A(z) \sim \sum A_{i} z^{-i}$ in $S$ as $z \rightarrow \infty$.

Suppose $A_{0}$ has eigenvalues divided into two disjoint sequences. Then exists $R_{1} \geq R$ and actual solutions $G(z), B(z)$ to

$$
z^{1-r} G^{\prime}+G B=A(z) G
$$

analytic for $|z| \geq R_{1}$ in $S$ with opening $<\frac{\pi}{r}$.

Remark. Sibuya proved that the result holds for $S$ with opening $\frac{\pi}{r}+\varepsilon$ for $\varepsilon>0$.

Now suppose $A_{0}$ has pairwise distinct eigenvalues $\lambda_{1}, \ldots, \lambda_{n}$. Then $G_{0}^{-1} A_{0} G_{0}=$ $\Lambda=\operatorname{diag}\left(\lambda_{1}, \ldots, \lambda_{n}\right), B(z)=\operatorname{diag}\left(b_{1}(z), \ldots, b_{n}(z)\right)$ is diagonal, $b_{k}(z) \sim \sum b_{\ell}^{(k)} z^{-\ell}$ with $b_{0}^{(k)}=\lambda_{k}$ in some $S$. Then the sysmte

$$
\frac{d \tilde{Y}}{d z}=z^{r-1} B(z) \tilde{Y}
$$

so exists a fundamental matrix solution $\tilde{Y}(z)=\operatorname{diag}\left(y_{1}(z), \ldots, y_{m}(z)\right)$ : by separation of variables

$$
y(z)=\exp \left(\int_{a}^{z} t^{r-1} b(t) d t\right)
$$

where the integration should be taken on a path that lies in $S$.

$$
\begin{aligned}
\int_{a}^{z} t^{r-1} b(t) d t & =\int_{a}^{z} t^{r-1}\left(b_{0}+\frac{b_{1}}{t}+\cdots+\frac{b_{n}}{t^{n}}\right) d t+\int_{a}^{z} \underbrace{t^{r-1}\left(b(t)-\sum_{j=0}^{r} \frac{b_{j}}{t^{j}}\right.}_{\bullet}) d t \\
& =q(z)+b_{n}(z)+\mathrm{const}+\int_{a}^{z} t^{r-1}\left(b(t)-\sum_{j=0}^{r} \frac{b_{j}}{t^{j}}\right) d t
\end{aligned}
$$

The term $(\bullet)$ vanishes in $S$ as $\frac{1}{t^{2}}$. Thus the limit of the integral as $z \rightarrow \infty$ exists, so

$$
\int_{a}^{z} t^{r-1}\left(b(t)-\sum_{j=0}^{r} \frac{b_{j}}{t^{j}}\right) d t=C+\underbrace{\int_{\infty}^{z} t^{r-1}\left(b(t)-\sum_{j=0}^{r} \frac{b_{j}}{t^{j}}\right)}_{h(z)} d t
$$

and

$$
h(z) \sim \int_{\infty}^{z} \sum_{\ell=1}^{\infty} \frac{b_{\ell+r}}{t^{\ell+1}} d t=\sum_{\ell=1}^{\infty} \int_{\infty}^{z} \frac{b_{\ell+r}}{t^{\ell+1}} d t=\sum_{\ell=1}^{\infty} \frac{b_{\ell+r}}{-\ell z^{\ell}}
$$

Thus in conclusion,

$$
y(z)=\text { const } \cdot z^{b_{r}} e^{q(z)} e^{h(z)}=f(z) z^{b_{r}} e^{q(z)}
$$

wher $f(z)$ is analytic in $S$. We can choose $f(z) \sim 1+\sum_{k=1}^{\infty} f_{k} z^{-k}$ and

$$
Y(z)=G_{0} \underbrace{G(z)\left(\begin{array}{ccc}
f_{1}(z) & & \\
& \ddots & \\
& & f_{n}(z)
\end{array}\right)}_{\mathcal{G}(z)} z^{B_{r}} e^{Q(z)}
$$

where

$$
Q(z)=\frac{B_{0}}{r} z^{r}+\frac{B_{1}}{r-1}+\cdots+B_{r-1} z .
$$

$\mathcal{G}(z) \sim F(z)=I+\sum_{k=1}^{\infty} \frac{F_{k}}{z^{k}}$ as $z \rightarrow 0$.

Definition (formal fundamental solution). $Y_{F}(z)=G_{0} F(z) z^{B_{r}} e^{Q(z)}$ is called the formal fundamental solution.

Remark. $\mathcal{G}(z) \sim F(z)$ asymptotically so $Y(z) z^{-B_{r}} e^{-Q(z)} \sim G_{0} F(z)$. By abuse of notation we will write $Y(z) \sim Y_{F}(z)$.

## Proposition 9.3.

- Suppose $\Lambda$ has been fixed. Then $B_{0}=\Lambda, B_{1}, \ldots, B_{r}$ are uniquely determined.
- Suppose $\Lambda$ and $G_{0}$ have been fixed. Then $Y_{F}(z)$ is uniquely determined.

Proof. Any formal solution has the form $Y_{F}(z) \cdot C$. Then

$$
\begin{aligned}
Y_{F}(z) C & =G_{0}\left(I+\sum_{k=1}^{\infty} F_{k} z^{-k}\right) z^{B_{r}} e^{Q(z)} C \\
& =G_{0}\left(I+\sum_{k=1}^{\infty} F_{k} z^{-k}\right) C z^{C^{-1} B_{r} C} e^{C^{-1} Q(z) C} \\
& =G_{0} C\left(I+\sum_{k=1}^{\infty} C^{-1} F_{k} C z^{-k}\right) z^{C^{-1} B_{r} C} e^{C^{-1} Q(z) C}
\end{aligned}
$$

We are requiring that $C^{-1} \Lambda C=\Lambda$. Since $\Lambda$ has distinct eigenvalues, this is equivalent to $C$ being diagonal so $C^{-1} B_{j} C=B_{j}$ for $j=1, \ldots, r$. This proves the first statement.

If in addition $G_{0}$ is fixed, $G_{0} C=G_{0}$ so $C=I$.
Now we adpot a different notation. Let $B_{1}=\Lambda_{1}, \ldots, B_{r}=\Lambda_{r}$. Then

$$
z^{B_{r}} e^{Q(z)}=\exp \left(\frac{\Lambda_{r}}{r} z^{r}+\cdots+\Lambda_{r-1} z+\Lambda_{r} \ln z\right)=: e^{\Lambda(z)}
$$

Remark. Though $Y_{F}$ is unique, this is not the case for $Y(z)$. In Wasow, the sector opening $<\frac{\pi}{r}$. If we increase the opening we can prove uniqueness.

Proposition 9.4. Two systems

$$
\frac{d Y}{d z}=z^{r_{1}-1} A_{1}(z) Y, \frac{d X}{d z}=z^{r_{2}-1} A_{2}(z) X
$$

are formally holomorphically equivalent if and only if $\pi_{1}=\pi_{2}$ and they have the same $\Lambda, \Lambda_{1}, \ldots, \Lambda_{r}$ where $\Lambda$ is a digonal form of both $A_{0}^{(1)}$ and $A_{0}^{(2)}$. These are called formal invariants.

Proof. Suppose they have the same formal invariants. Then

$$
\begin{aligned}
& Y_{F}(z)=G_{0}^{(1)} F_{1}(z) \exp \left(\frac{\Lambda}{r_{1}} z^{r_{1}}+\cdots+\Lambda_{r_{1}} \ln z\right) \\
& X_{F}(z)=G_{0}^{(2)} F_{2}(z) \exp \left(\frac{\Lambda}{r_{1}} z^{r_{1}}+\cdots+\Lambda_{r_{1}} \ln z\right)
\end{aligned}
$$

Then $X_{F}(z)=H(z) Y_{F}(z)$ where $H(z)=G_{0}^{(2)} F_{2}(z)\left(G_{0}^{(1)} F_{1}(z)\right)^{-1}$.
Conversely suppose $X(z)=H(z) Y(z)$ wher $\operatorname{det} H_{0} \neq 0$. Let $\Lambda$ be a diagonal form of $A_{0}^{(1)}$. Then exists

$$
Y_{F}(z)=G_{0}^{(1)} F_{1}(z) \exp \left(\frac{\Lambda}{r_{1}} z^{r_{1}}+\cdots+\Lambda_{r_{1}} \ln z\right)
$$

so by assumption

$$
X_{F}(z)=\underbrace{H_{z}(z) G_{0}^{(1)} F_{1}(z)}_{\mathcal{F}(z)} \exp \left(\frac{\Lambda}{r_{1}} z^{r_{1}}+\cdots+\Lambda_{r_{1}} \ln z\right) .
$$

The equation for $X_{F}$ is

$$
\frac{d X_{F}}{d z}=\left(\frac{d \mathcal{F}}{d z} \mathcal{F}^{-1}+\mathcal{F}\left(\Lambda z^{r_{1}-1}+\cdots+\frac{\Lambda_{r_{1}}}{z}\right) \mathcal{F}^{-1}\right) X_{F}=z^{r_{1}-1}(\underbrace{\mathcal{F}_{0} \Lambda \mathcal{F}_{0}^{-1}}_{A_{0}^{(2)}}+\sum \frac{A_{i}^{(2)}}{z^{i}}) X_{F}
$$

so $A_{0}^{(2)}$ has diagonal form $\Lambda, r_{2}=r_{1}$ and also $\Lambda_{1}, \ldots, \Lambda_{r}$ are the same.

### 9.2 Case of non-distinct eigenvalues

Theorem 9.5. If $A(z)$ is holomorphic in $S, A(z) \sim \sum A_{i} z^{i}$ then for every "sufficiently small" subsector of $S$ there is fundamental matrix solution

$$
Y(z)=G(z) z^{L} e^{Q(z)}
$$

where $Q(z)$ is diagonal polynomial in $z^{1 / N}$ for some natural number $N, L$ is inte general not diagonal, $G(z) \sim z^{k_{0} / N} \sum_{j=0}^{\infty} \frac{F_{j}}{z^{j / N}}$ in some small sector for $k_{0} \in \mathbb{Z}$. In general $\operatorname{det} F_{0}=0 . z=\infty$ is called a ramified singularity.

Example: Airy equation The Airy equation is the equation

$$
u^{\prime \prime}=z u .
$$

By a substitution

$$
\frac{d y}{d z}=\left(\begin{array}{ll}
0 & 1 \\
z & 0
\end{array}\right) y=z\left(\left(\begin{array}{ll}
0 & 0 \\
1 & 0
\end{array}\right)+\frac{1}{z}\left(\begin{array}{ll}
0 & 1 \\
0 & 0
\end{array}\right)\right) y .
$$

By the gauge transformation $y=\left(\begin{array}{cc}0 & 1 \\ 1 & 0\end{array}\right) y_{1}$, we write

$$
\frac{d y_{1}}{d z}=z\left(\left(\begin{array}{ll}
0 & 1 \\
0 & 0
\end{array}\right)+\frac{1}{z}\left(\begin{array}{ll}
0 & 0 \\
1 & 0
\end{array}\right)\right) y_{1}
$$

from which we see that there are two identical eigenvalues $\lambda_{1}=\lambda_{2}=0$.
If we apply the shearing transformation $y_{1}=\left(\begin{array}{cc}1 & 0 \\ 0 & z^{-1 / 2}\end{array}\right) y_{2}$ so

$$
\frac{d y_{2}}{d z}=z^{1 / 2}\left(\left(\begin{array}{ll}
0 & 1 \\
1 & 0
\end{array}\right)+\frac{1}{2 z^{3 / 2}}\left(\begin{array}{ll}
0 & 0 \\
0 & 1
\end{array}\right)\right) y_{2}
$$

so at the cost of a fractional power we now have two distinct eigenvalues. Let $t=c z^{1 / 2}$ so $\frac{d}{d z}=\frac{c^{2}}{2} \frac{1}{t} \frac{d}{d t}$. Set $c=2^{1 / 3}$,

$$
\frac{d Y_{2}}{d t}=t^{2}\left(\left(\begin{array}{ll}
0 & 1 \\
1 & 0
\end{array}\right)+\frac{1}{t^{3}}\left(\begin{array}{ll}
0 & 0 \\
0 & 1
\end{array}\right)\right) y_{2}
$$

This has Poincaré rank 3. Eigenvalues of the term in the parethesis is

$$
\pm 1+\frac{1}{2 t^{3}}+O\left(\frac{1}{t^{6}}\right)
$$

so

$$
Y_{2}(z)=G_{0} G(t) t^{\Lambda_{3}} \exp \left(\frac{\Lambda}{3} t^{3}+\frac{\Lambda_{1}}{2 t}\right)
$$

SO

$$
Y(z)=\left(\begin{array}{cc}
0 & 1 \\
1 & 0
\end{array}\right)\left(\begin{array}{cc}
1 & 0 \\
0 & z^{-1 / 2}
\end{array}\right)\left(\begin{array}{cc}
-\gamma_{1} & \gamma_{2} \\
\gamma_{1} & \gamma_{2}
\end{array}\right) G\left(z^{1 / 2}\right) z^{\frac{1}{4} I} \exp \left(-\frac{2}{3} z^{3 / 2}\left(\begin{array}{cc}
1 & 0 \\
0 & -1
\end{array}\right)\right)
$$

### 9.3 Stokes phenomenon

Now back to the case $A_{0}$ with distinct eigenvalues. Suppose $Y(z) \sim Y_{F}(z)$ in $S$. We can extend asympototics to $S \supseteq S^{\prime}$ but in general we meet "separating rays" beyond which the asymptotics no longer holds. This is called Stokes phenomenon. The ray is called Stokes ray.

Exercise. As $z \rightarrow \infty$,

$$
f(z)=1+e^{z}= \begin{cases}\sim 1 & \operatorname{Re} z<0 \\ \text { oscilates } & \operatorname{Re} z=0 \\ \infty & \operatorname{Re} z>0\end{cases}
$$

Consider the $\lambda$-plane. Choose $\eta \in \mathbb{R}$ such that $\arg \left(\lambda_{j}-\lambda_{k}\right) \neq \eta(\bmod \pi)$. This is called an admissible direction in the $\lambda$-plane. Take any determination of $\arg \left(\lambda_{j}-\lambda_{k}\right)$, for example $\eta-2 \pi<\arg \left(\lambda_{j}-\lambda_{k}\right)<\eta$ for $j \neq k$.

Definition (Stokes ray). The Stokes rays associated with $\left(\lambda_{j}, \lambda_{k}\right)_{j \neq k}$ are infinitely many rays in $\widetilde{\mathbb{C} \backslash\{0\}}$ such that

$$
\operatorname{Re}\left(\left(\lambda_{j}-\lambda_{k}\right) z^{n}\right)=0, \operatorname{im}\left(\left(\lambda_{j}-\lambda_{k}\right) z^{n}\right)<0
$$

Proposition 9.6. To $\left(\lambda_{j}, \lambda_{k}\right)$ the associated rays are

$$
\arg z=\theta_{j k}+2 \pi \frac{N}{r}
$$

for all $N \in \mathbb{Z}, \theta_{j k}=\frac{1}{r}\left(\frac{3 \pi}{2}-\arg \left(\lambda_{j}-\lambda_{k}\right)\right)$.
Proof. Exercise.

There is an even number $2 \mu$ of possible values of $\arg \left(\lambda_{j}-\lambda_{k}\right)$. Label $\arg \left(\lambda_{j}-\right.$ $\lambda_{k}$ ) with $\eta_{\tilde{\nu}}$ for $0 \leq \tilde{\nu} \leq 2 \mu-1$. If $\eta$ is an admissible direction in $\lambda$-plane,

$$
\eta>\eta_{0}>\eta_{1}>\ldots \eta-\pi>\eta_{\mu}>\cdots>\cdot_{2 \pi-1}>\eta-2 \pi .
$$

Let $\pi=\frac{3 \pi}{2}-\eta$. Stokes rays have directions $\tau_{\tilde{\nu}}=\left(\frac{3 \pi}{2}-\eta_{\tilde{\nu}}\right) \frac{1}{r}$ and they satisfy

$$
\tau<\tau_{0}<\tau_{1}<\cdots<\tau_{\mu-1}<\tau+\frac{\pi}{2}<\tau_{\mu}<\cdots<\tau_{2 \pi-1}<\tau+\frac{2 \pi}{r}
$$

$\tau$ is called admissible direction in the z-plane.
From the proposition all Stokes rays are $\tau_{\nu}=\tau_{\tilde{\nu}}+\frac{2 N \pi}{r}$. In fact, we can express all Stokes rays by

$$
\tau_{\tilde{\nu}+N \mu}:=\tau_{\tilde{\nu}}+N \frac{\pi}{r}
$$

for $0 \leq \tilde{\nu} \leq \mu-1$.
One can also conclude that a section not containing a Stokes ray has opening $<\frac{\pi}{r}$.

Lemma 9.7. Suppose $Y(z)$ is a fundamental matrix solution such that $Y(z) \tilde{Y}_{F}(z)$ as $z \rightarrow \infty$ in $S$. Assume exists $\tilde{S}$ that does not contain Stokes rays and $S \cap \tilde{S} \neq \emptyset$, then $Y(z) \sim Y_{F}(z)$ in $S \cup \tilde{S}$.

Proof. $\tilde{S}$ has central opening $<\frac{\pi}{r}$. By Wasow exists $\left.\tilde{( } z\right) \sim Y_{F}(z)$ in $\tilde{S}$. Then $Y(z)=\tilde{Y}(z) C$ for some $C$. As $Y(z)=G_{0} \mathcal{G}(z) e^{\Lambda(z)}, \tilde{Y}(z)=G_{0} \tilde{\mathcal{G}}(z) e^{\Lambda(z)}$ so

$$
e^{\Lambda(z)} C e^{-\Lambda(z)}=\tilde{\mathcal{G}}(z)^{-1} \mathcal{G}(z)
$$

but $\mathcal{G}$ and $\tilde{\mathcal{G}}$ has exactly the same same asymptotic expansion in $S \cap \tilde{S}$ so RHS $\sim I$ as $z \rightarrow \infty$. Then

$$
\begin{equation*}
e^{\Lambda_{i}(z)-\Lambda_{j}(z)} C_{i j}=\exp \left(\frac{\lambda_{i}-\lambda_{j}}{r} z^{r}\left(1+O\left(\frac{1}{z}\right)\right)\right) C_{i j} \sim \delta_{i j} \tag{*}
\end{equation*}
$$

By hypothesis in $S \cap \tilde{S}$ there are no Stokes rays so the sign of $\operatorname{Re}\left(\lambda_{i}-\lambda_{j}\right) z^{r}$ is fixed in $S \cap \tilde{S}$, and also in $\tilde{S}$. Thus $(*)$ holds in $\tilde{S}$ and

$$
\mathcal{G}(z)=\tilde{\mathcal{G}}(z) \underbrace{e^{\Lambda(z)} C e^{-\Lambda(z)}}_{\sim I}
$$

so the asymptotic expansion holds in $S \cup \tilde{S}$.
As a corollary we have

Theorem 9.8 (extension). Let $Y(z)$ be a fundamental matrix solution which by Wasow's result satisfies $Y(z) \sim Y_{F}(z)$ in an $S$ containing a set of basic Stokes rays. Then $Y(z) \sim Y_{F}(z)$ in an open sector $\mathcal{S}$ containing $S$. $\mathcal{S}$ extending up to the nearest Stokes rays outside $S$.

Theorem 9.9 (uniqueness). Let $Y(z)$ be a fundamental matrix solution such that $Y(z) \sim Y_{F}(z)$ in a sector containing a set of basic Stokes rays. Then $Y(z)$ is unique.

Proof. Suppose exists $\tilde{Y}(z) \sim Y_{F}(z)$ in the sector. Then as usual $Y(z)=\tilde{Y}(z) C$ and

$$
e^{\Lambda(z)} C e^{-\Lambda(z)}=\tilde{\mathcal{G}}(z)^{-1} \mathcal{G}(z) \sim I
$$

so

$$
\begin{equation*}
e^{\Lambda_{i}(z)-\Lambda_{j}(z)} C_{i j} \sim \delta_{i j} . \tag{*}
\end{equation*}
$$

Since the sector contains basic rays, $\operatorname{Re}\left(\lambda_{i}-\lambda_{j}\right) z^{n}$ changes signas $z$ varies in the sector. Thus there exists a subsector where $e^{\left(\lambda_{i}-\lambda_{j}\right) z^{r}} \rightarrow \infty$ as $z \rightarrow \infty$, so $(*)$ holds if and only if $C_{i j}=\delta_{i j}$ so $C=I$.

Remark. We can take the sector to be $\mathcal{S}$ in the extension theorem.

### 9.4 Stokes matrix

Let $Y_{1}(z), Y_{2}(z)$ are fundamental solutions such that $Y_{j}(z) \sim Y_{F}(z)$ in $S_{j}$ such that the opening of $S_{j}>\frac{\pi}{r}$. Suppose $S_{1} \cap S_{2} \neq \emptyset$ and does not contain Stokes rays.

Definition (Stokes matrix). The Stokes matrix is the connection matrix $S$ defined by

$$
Y_{2}(z)=Y_{1}(z) S
$$

for $z \in S_{1} \cap S_{2}$ (and analytically exnteded in $z$ ).
Exercise. Consider Stokes rays associated with $\left(\lambda_{j}, \lambda_{k}\right)$. Prove that for $\arg z=$ $\theta_{j k}+\frac{2 \pi N}{r}+\delta, \delta \in \mathbb{R}$ (taken modulo $\frac{2 \pi}{r}$ ),

$$
\operatorname{Re}\left(\lambda_{j}-\lambda_{k}\right) z^{n} \begin{cases}<0 & -\frac{\pi}{r}<\delta<0 \\ =0 & \delta=0, \pm \frac{\pi}{r} \\ >0 & 0<\delta<\frac{\pi}{r}\end{cases}
$$

Definition. If in an open sector $\operatorname{Re}\left(\lambda_{j}-\lambda_{k}\right) z^{n}>0$ we write $\lambda_{j} \succ \lambda_{k}$ and say $\lambda_{j}$ is dominant. If $\operatorname{Re}\left(\lambda_{j}-\lambda_{k}\right) z^{n}<0$ then write $\lambda_{j} \prec \lambda_{k}$.

If $S$ does not contain Stokes rays then $\prec$ defines an ordering in $\left\{\lambda_{1}, \ldots, \lambda_{n}\right\}$.

Proposition 9.10. Suppose $Y_{2}=Y_{1} S$ for some Stokes matrix $S$ in $S_{1} \cap S_{2}$. Then $S$ has the following form: $S_{j j}=1$ for all $j, S_{j k}=0$ for $\lambda_{j} \succ \lambda_{k}$ in $S_{1} \cap S_{2}$,

Proof. As usual write $Y_{a}(z)=G_{0} \mathcal{G}_{a}(z) e^{\Lambda(z)}$. Then

$$
e^{\Lambda_{j}(z)-\Lambda_{k}(z)} S_{j k} \sim \delta_{j k}
$$

This holds if and only if $S_{j j}=1$ and $S_{j k}=0$ for $\operatorname{Re}\left(\lambda_{j}-\lambda_{k}\right) z^{n} \geq 0$, which is precisely the definition of $\lambda_{j} \succ \lambda_{k}\left(S_{j k}\right.$ can be any number for $\operatorname{Re}\left(\lambda_{j}-\lambda_{k}\right) z^{n}<$ $0)$.

Let $\mathcal{S}_{\nu}=S\left(\tau_{\nu-\mu}, \tau_{\nu+1}\right)$ be the open sector containing basic rays $\tau_{\nu-\mu+1}, \ldots, \tau_{\nu}$. Let $\mathcal{S}_{\nu+\mu}=S\left(\tau_{\nu}, \tau_{\nu+\mu+1}\right)$ and so on. On each sector $\mathcal{S}_{\nu+h \mu}, h \in \mathbb{Z}$, exists a unique solution $Y_{\nu+h \mu(z)} \sim Y_{F}(z)$. There are Stokes matrices relating overlapping sector, which we label using

$$
Y_{\nu+(h+1) \mu}(z)=Y_{\nu+h \mu} S_{\mu+h \mu} .
$$

Exercise. Show that $\left(S_{\nu}\right)_{j k}=0$ if and only if $\left(S_{\nu+\mu}\right)_{j k}=0$, if and only if $\left(S_{\nu-\mu}\right)_{k j}=0$ (so they are upper/lower triangular alternating in $\nu$ ).

### 9.5 Monodromy

Consider the formal solution $Y_{F}(z)=G_{0} F(z) e^{\Lambda(z)}$. As $\Lambda(z)$ can be written as a polynomial plus $\Lambda_{r} \log z$,

$$
Y_{F}\left(z e^{2 \pi i}\right)=G_{0} F(z) e^{\Lambda(z)} e^{2 \pi i \Lambda_{r}} .
$$

We call $e^{2 \pi i \Lambda_{r}}$ formal monodromy.
Given $z \in \mathcal{S}_{\nu}, z_{\nu} \mapsto z_{\nu+2 r \mu}=z_{\nu} \cdot e^{2 \pi i} \in \mathcal{S}_{\nu+2 r \mu}$.

## Theorem 9.11.

1. $Y_{\nu+2 r \mu}\left(z_{\nu+2 \pi \mu}\right)=Y_{\nu}\left(z_{\nu}\right) \cdot e^{2 \pi i \Lambda_{r}}$.
2. $Y_{\nu+2 r \mu}(z)=Y_{\nu}(z) S_{\nu} S_{\nu+\mu} \cdots S_{\nu+(2 r-1) \mu}$.
3. $Y_{\nu}\left(z e^{2 \pi i}\right)=Y_{\nu}(z) e^{2 \pi i \Lambda_{r}}\left(S_{\nu} \cdots S_{\nu+(2 r-1) \mu}\right)^{-1}=Y_{\nu}(z)\left[M_{\infty}^{(\nu)}\right]^{-1}$.

Proof. 1 follows from formal monodromy and uniqueness of asymptotics. 2 is repeated application of Stokes matrix. Combining them gives 3.

Proposition 9.12. For every $\nu \in \mathbb{Z}$, we have the relation

$$
S_{\nu+2 r \mu}=e^{-2 \pi i \Lambda_{r}} S_{\nu} e^{2 \pi i \Lambda_{r}} .
$$

Proof.

$$
\begin{aligned}
Y_{\nu+(2 r+1) \mu}\left(z_{\nu+(2 r+1) \mu}\right) & =Y_{\nu+\mu}\left(z_{\nu+\mu}\right) e^{2 \pi i \Lambda_{r}} \\
& =Y_{\nu}\left(z_{\nu}\right) S_{\nu} \cdot e^{2 \pi i \Lambda_{r}}
\end{aligned}
$$

LHS can also be expressed as

$$
Y_{\nu+2 r \mu}\left(z_{\nu+(2 r+1) \mu}\right) S_{\nu+2 r \mu}=Y_{\nu}\left(z_{\nu}\right) e^{2 \pi i \Lambda_{r}} S_{\nu+2 r \mu} .
$$

## Theorem 9.13.

1. $\Lambda_{r}, S_{\nu}, S_{\nu+\mu}, \ldots, S_{n+(2 r-1) \mu}$ generated $S_{\nu+h \mu}$ for all $h \in \mathbb{Z}$.
2. They are sufficient to compute the monodromy at $\infty$ of every $Y_{\nu+h \mu}$.

Definition (complete set of Stokes matrices). $S_{\nu}, \ldots, S_{\nu+(2 r-1) \mu}$ is called a complete set of Stokes matrices. $S_{\nu}, \ldots, S_{n+(2 r-1) \mu}, \Lambda_{r}$ are called monodromy data at $z=\infty$.

Theorem 9.14. $\frac{d Y}{d z}=z^{r-1} A(z) Y, \frac{d \tilde{Y}}{d z}=z^{\tilde{r}-1} \tilde{A}(z) \tilde{Y}$ are holomorphically equivalent if and only if $=\tilde{r}$, have the same $\Lambda, \Lambda_{1}, \ldots, \Lambda_{r}$, and for some $\nu_{0}$ they have the same $S_{\nu_{0}}, \ldots, S_{\nu_{0}+(2 r-1) \mu}$ (this is true for all $\nu$ ).

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